

1     **The Role of a Polar/Subtropical Jet Superposition in the May**

2                     **2010 Nashville Flood**

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## ABSTRACT

Contributions to the increased poleward moisture flux that characterized the second day of the May 1-3 Nashville Flood of 2010 are examined from the perspective of polar and subtropical jet superposition and its influence on the secondary ageostrophic circulation. Employing the Sawyer-Eliassen circulation equation, the analysis reveals that the poleward moisture flux attributed to the jet increased nearly 120% prior to the second day of the event in response to the superposed jet's ageostrophic circulation, helping to further fuel the production of heavy rainfall.

The full Sawyer-Eliassen circulation associated with the superposed jet is further partitioned into its geostrophic and diabatic components. The geostrophic forcing drove middle tropospheric ascent that fueled the production of deep convection and the record rainfall. The diabatic component, through forcing lower tropospheric ascent and vigorous lower tropospheric poleward moisture flux, provided the link between the tropical moisture and the deep convective environment. Since superposed jets, by their nature, develop on the poleward edge of the tropical or subtropical air, it is suggested that such a mutually reinforcing interaction between these two component forcings of the secondary circulation may routinely characterize the involvement of superposed jet structures in high impact weather events.

## 21 **1. Introduction**

22 During the first two days of May 2010, two consecutive mesoscale convective systems  
23 (MCSs) were responsible for historic rainfall accumulations in excess of 180 mm (7 in.) over  
24 a large portion of Tennessee, southern Kentucky, and northern Mississippi (Fig. 1). A few  
25 locations saw significantly higher rainfall totals, according to the National Weather Ser-  
26 vice office in Nashville, with Camden, Tennessee tallying 493 mm (19.41 in.) and Nashville  
27 recording 344 mm (13.54 in.) during that two-day period. In addition to the heavy rainfall,  
28 both days were characterized by tornado outbreaks, as the environment was weakly strati-  
29 fied and favorable for severe convective development. This combination of sensible weather  
30 had enormous and wide-ranging impacts across the entire region, closing numerous roads,  
31 resulting in 26 flood related fatalities, causing around \$2 billion in property damage in the  
32 greater Nashville area, and swelling area rivers to record crests. Specifically, the Cumberland  
33 River at Nashville recorded a crest of 15.8 m (51.9 ft), which was 1.3 m (4.3 ft) higher than  
34 the previous record at the station in the post flood control era (National Weather Service  
35 2011; hereafter NWS 2011).

36 One of the most notable aspects of this event, as diagnosed in the work of Moore et al.  
37 (2012, hereafter M12), was the presence of an anomalous and narrow plume of enhanced  
38 water vapor transport, or an atmospheric river (Newell et al. 1992 and Zhu and Newell 1998),  
39 that extended from the Gulf of Mexico northward into the eastern United States. While  
40 several recent studies have investigated the impacts of atmospheric rivers on orographically-  
41 forced precipitation events (e.g. Ralph et al. 2006; Stohl et al. 2008; Guan et al. 2010), the  
42 M12 study demonstrated that atmospheric rivers can also play a considerable role in heavy  
43 rainfall events that are synoptically forced, such as those that occur over the central and  
44 eastern United States.

45 Specifically, M12 found that a southerly low-level jet, driven by a strong geopotential  
46 height gradient between a lee trough along the east coast of Mexico and a strong subtropical  
47 ridge north of the Caribbean Sea, facilitated much of the anomalous moisture transport out

48 of Central America and into the southern Mississippi River Valley. Furthermore, they noted  
49 that this moisture transport strengthened over the northern Gulf of Mexico and southern  
50 Mississippi River Valley in the hours preceding the second day of the event (their Fig. 6).  
51 This finding was also noted in subsequent studies of the Nashville flood (Durkee et al. 2012;  
52 Lackmann 2013; Lynch and Schumacher 2013). These studies have demonstrated that this  
53 persistent and increased moisture transport into the region over the two-day period, in  
54 conjunction with ascent along stationary, convectively generated outflow boundaries, aided  
55 in the production of heavy rainfall across portions of Tennessee, Kentucky, and northern  
56 Mississippi.

57 Coincident with the increase in moisture transport prior to the second day of heavy  
58 rainfall, however, was a relatively rare vertical superposition of the normally distinct polar  
59 and subtropical jet streams and an attendant acceleration of jet wind speeds (see M12, their  
60 Fig. 4). Observational analysis by Defant and Taba (1957, hereafter DT57) of tropopause  
61 temperature (their Fig. 3) demonstrates that, in such a superposition, the upper tropospheric  
62 and lower stratospheric baroclinicity associated with each jet is intensified. As a result, a  
63 superposed jet structure possesses an anomalous fraction of the pole-to-equator temperature  
64 gradient (manifest as available potential energy (APE)). This suggests that much stronger  
65 upper tropospheric and lower stratospheric fronts and an anomalously deep layer of vertical  
66 shear, as required by the increased horizontal baroclinicity, accompany the relatively rare  
67 superposition of the polar and subtropical jets.

68 The development of intensified frontal structure associated with the superposed jet is  
69 often attended by a strengthening of its transverse, ageostrophic secondary circulation, di-  
70 agnosable using the Sawyer-Eliassen circulation equation (Sawyer 1956; Eliassen 1962). Such  
71 ageostrophic circulations have been shown in numerous studies to play an important role in  
72 the production of sensible weather. For example, much attention has been focused on upper  
73 tropospheric fronts, which can form as a result of the differential vertical motions associated  
74 with Sawyer-Eliassen circulations and are an important part of the extratropical cyclone life

75 cycle (e.g. Uccellini et al. 1985; Whitaker et al. 1988; Barnes and Colman 1993; Lackmann  
76 et al. 1997). Additionally, the circulations associated with upper tropospheric fronts have  
77 been shown to play an important role in the development of convective precipitation events,  
78 as first suggested by Omoto (1965) and further demonstrated by Hobbs et al. (1990) and  
79 Martin et al. (1993). While a number of studies have qualitatively considered the moisture  
80 flux accomplished by the lower tropospheric horizontal branches of ageostrophic jet circula-  
81 tions (e.g. Uccellini and Johnson 1979; Uccellini et al. 1984; Uccellini and Kocin 1987), direct  
82 quantification of these effects has not received as much attention in the literature. Further-  
83 more, if the static stability is low in a given region, as it was over the southern Mississippi  
84 River Valley on 1-2 May 2010, a Sawyer-Eliassen circulation can occupy a considerable depth  
85 of the troposphere. In such a situation, the horizontal winds associated with the secondary  
86 circulation near the surface are capable of significant contributions to the moisture transport  
87 into the region.

88 While M12 and Durkee et al. (2012) acknowledge a strengthening of both the jet and  
89 moisture flux prior to the second day of heavy rainfall, they do not investigate the link  
90 between these processes. Consequently, the modulation of the structure and intensity of the  
91 Sawyer-Eliassen circulation by the diabatic residue of the heavy rainfall that characterized  
92 this event remains to be considered. In order to address these issues, the present study  
93 aims to 1) quantify the contribution to the poleward moisture flux made by the superposed  
94 jet's ageostrophic circulation and 2) examine the impact that both geostrophic and diabatic  
95 forcing may have had in determining the strength and sense of the overall ageostrophic  
96 circulation.

97 The remainder of this study is organized as follows. Section 2 gives an overview of  
98 the methodology used to identify superposed jets as well as background on the Sawyer-  
99 Eliassen circulation equation. Section 3 provides a brief synoptic overview of the flooding  
100 event. Section 4 discusses the impacts of the Sawyer-Eliassen circulations during each day of  
101 the event and further dissects the forcing responsible for the superposed jet's ageostrophic

102 circulation. Finally, Section 5 presents a discussion and conclusions.

## 103 2. Methodology

104 This study is performed using model analyses from the National Centers for Environmen-  
105 tal Prediction (NCEP) Global Forecast System (GFS) at 6-h intervals with a horizontal grid  
106 spacing of  $1.0^\circ \times 1.0^\circ$  and a vertical grid spacing of 50 hPa (25 hPa between 1000 hPa and  
107 900 hPa). In order to accommodate the identification scheme that follows, these data were  
108 bi-linearly interpolated onto isentropic surfaces at 5K intervals from 300K to 370K using  
109 programs within the General Meteorology Package (GEMPAK) (desJardins et al. 1991).

### 110 a. *Jet Identification*

111 The work by DT57 identified the characteristic three-step pole-to-equator tropopause  
112 structure shown in Fig. 2 (modified from their Fig. 13), wherein each step is separated from  
113 its neighbors by the presence of a westerly wind maximum. They found that, on average, the  
114 tropical tropopause<sup>1</sup> was found at around 90 hPa (17-18 km) and extended to 30°N, roughly  
115 the poleward edge of the Hadley Cell. Near that latitude the tropopause height abruptly  
116 lowers to about 200 hPa (12 km), with the subtropical jet nestled within that break in  
117 the tropopause (e.g. Loewe and Radok 1950; Yeh 1950; Koteswaram 1953; Mohri 1953;  
118 Koteswaram and Parthasarathy 1954; Sutcliffe and Bannon 1954; Krishnamurti 1961; Riehl  
119 1962). Poleward of this feature was what DT57 termed the “middle tropopause” located  
120 around 250 hPa. The polar jet is found in the break between the middle tropopause and the  
121 even lower polar tropopause (300 hPa) near 50°N. While relatively modest baroclinicity in  
122 the upper troposphere and lower stratosphere characterizes the subtropical jet, the polar jet

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<sup>1</sup>DT57 identified the tropopause via analysis of soundings. The tropopause was identified at the elevation of “a noticeable change of tropospheric lapse rate to an isothermal layer or to an increase of temperature with height.” p. 261, DT57.

123 sits atop the strongly baroclinic, tropospheric-deep polar front (e.g. Namias and Clapp 1949;  
124 Palmén and Newton 1948; Newton 1954; Palmén and Newton 1969; Keyser and Shapiro 1986;  
125 Shapiro and Keyser 1990).

126 The DT57 analysis also demonstrated the utility of maps of tropopause height (in hPa)  
127 for locating the position of the jets. On these maps, one of which is shown in Fig. 3 (mod-  
128 ified from DT57 Fig. 2), DT57 referred to sharp, isolated, and easily identifiable gradients  
129 of tropopause height as “breaklines”. These breaklines were found to be coincident with the  
130 respective jet maxima (e.g. the subtropical jet is found at the breakline between the tropical  
131 and middle tropopause). While such an analysis demonstrates that these jets typically oc-  
132 cupy different latitude bands, substantial meanders in the jets are common. Additionally, the  
133 characteristic latitudinal separation between the two structures occasionally disappears, as it  
134 does in Fig. 3 south of Iceland over the North Atlantic, where the polar and subtropical jets  
135 vertically superpose. These observations of the tropopause structure, both climatologically  
136 and instantaneously, form the theoretical basis for the following jet identification scheme.

137 The identification scheme for the polar, subtropical, and superposed jet streams is de-  
138 scribed with reference to the features illustrated in Fig. 4. Figure 4a depicts an example  
139 of clearly separate polar and subtropical jets in the eastern North Pacific. A vertical cross  
140 section through these distinct features unambiguously identifies the separate jet cores (Fig.  
141 4b). From this cross section, it is clear that the core of the polar jet, located at approximately  
142 300 hPa, is largely contained within the 315-330K isentropic layer while the subtropical jet  
143 core, located at approximately 200 hPa, occupies the 340-355K layer. Additionally, both  
144 the polar and the subtropical jets lie at the low potential vorticity (PV) edge of the strong  
145 horizontal PV gradient that separates the upper troposphere from the lower stratosphere in  
146 their respective layers. With these attributes in mind, the identification scheme evaluates  
147 characteristics of the PV and wind speed distributions in each grid column of analysis data.  
148 Within the 315-330K (340-355K) layer, whenever the magnitude of the PV gradient within

149 the 1-3 PVU channel exceeds an empirically determined threshold value<sup>2</sup> *and* the integrated  
150 wind speed in the 400-100 hPa layer exceeds 30 m s<sup>-1</sup>, a polar (subtropical) jet is identified  
151 in that grid column. The occurrence of both polar and subtropical jet characteristics in  
152 a single grid column identifies a jet superposition event at that time in that grid column.  
153 An example of a jet superposition event is shown in a plan view in Fig. 4c. Not until a  
154 vertical slice through the jet core is examined can the superposition be identified (Fig. 4d).  
155 Notice that, rather than the three-step tropopause structure identified by DT57 and shown  
156 in Fig. 4b, a superposed jet is characterized by a two-step tropopause structure with a steep  
157 tropopause wall from the polar to the tropical tropopause. This nearly vertical PV wall  
158 (from roughly 550 hPa to 150 hPa in this case) is the leading structural characteristic of a  
159 superposed jet.

160 *b. Sawyer-Eliassen Circulation Equation*

161 A particularly useful way to interrogate the vertical circulations associated with jet-front  
162 structures, in nearly straight flow, is afforded by the Sawyer-Eliassen circulation equation  
163 (Sawyer 1956; Eliassen 1962):

$$(-\gamma \frac{\partial \theta}{\partial p}) \frac{\partial^2 \psi}{\partial y^2} + (2 \frac{\partial M}{\partial p}) \frac{\partial^2 \psi}{\partial p \partial y} + (-\frac{\partial M}{\partial y}) \frac{\partial^2 \psi}{\partial^2 p} = Q_g - \gamma \frac{\partial}{\partial y} (\frac{d\theta}{dt}) \quad (1)$$

164 where  $\gamma$  is a constant on isobaric surfaces ( $\gamma = (R/fp_0)(p_0/p)^{c_v/c_p}$ ),  $p_0 = 1000$  hPa,  $c_v = 718$   
165 J kg<sup>-1</sup> K<sup>-1</sup>,  $c_p = 1004$  J kg<sup>-1</sup> K<sup>-1</sup>,  $R$  is the gas constant for dry air,  $\theta$  is the potential tem-  
166 perature, and  $f$  is the Coriolis parameter.  $M$  is the absolute geostrophic momentum ( $M =$   
167  $U_g - fy$ ) and  $U_g$  and  $V_g$  are the along- and across-front geostrophic winds, respectively.  $Q_g$  is  
168 the geostrophic forcing term, which is the sum of the shearing ( $Q_{SH} = 2\gamma[(\partial U_g/\partial y)(\partial \theta/\partial x)]$ )  
169 and stretching deformation terms ( $Q_{ST} = 2\gamma[(\partial V_g/\partial y)(\partial \theta/\partial x)]$ ). The ageostrophic circula-  
170 tion lies in a plane transverse to the frontal boundary (jet axis) and is determined by the

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<sup>2</sup>The threshold values are  $1.4 \times 10^{-5}$  PVU m<sup>-1</sup> ( $1.4 \times 10^{-11}$  m K kg<sup>-1</sup> s<sup>-1</sup>) for the 315-330K layer and  $0.9 \times 10^{-5}$  PVU m<sup>-1</sup> for the 340-355K layer.

171 Sawyer-Eliassen streamfunction,  $\psi$ , such that  $v_{ag} = -\partial\psi/\partial p$  and  $\omega = dp/dt = \partial\psi/\partial y$ . Given  
172 the second-order nature of this differential equation, positive (negative) values for the forcing  
173 function correspond to negative (positive) values for the streamfunction and thermally direct  
174 (indirect) circulations. The coefficients of the second-order terms on the left hand side of  
175 (1) represent the static stability, baroclinicity, and inertial stability, respectively. For the  
176 full derivation and discussion of (1), the reader is referred to Eliassen (1962) or Keyser and  
177 Shapiro (1986).

178 From (1), it becomes evident that knowledge of the distribution of  $U_g$ ,  $V_g$ ,  $M$ ,  $\theta$ , and  
179  $d\theta/dt$ , in a particular case, allows for the calculation of the coefficients on the left hand side  
180 of (1) as well as the forcing function. Consequently, absolute temperature and geostrophic  
181 wind are extracted from each grid point in the GFS analysis at 50 hPa vertical intervals from  
182 1000 hPa to 50 hPa. These variables are then interpolated onto the selected vertical cross  
183 section perpendicular to the jet axis. Subsequently, all of the coefficients and geostrophic  
184 forcing terms in (1) are calculated from these interpolated variables at each grid point within  
185 the interior of the cross section. Model vertical motion and relative humidity data are also  
186 extracted from the GFS analysis and interpolated onto the cross sectional grid in order to  
187 determine  $d\theta/dt$ , or the rate of latent heating. Following the method of Emanuel et al.  
188 (1987), this term is calculated as:

$$\frac{d\theta}{dt} = \omega \left( \frac{\partial\theta}{\partial p} - \frac{\Gamma_m}{\Gamma_d} \frac{\theta}{\theta_e} \frac{\partial\theta_e}{\partial p} \right) \quad (2)$$

189 where  $\omega$  is the model vertical motion,  $\theta_e$  is the equivalent potential temperature, and  $\Gamma_m$  and  
190  $\Gamma_d$  are the moist and dry adiabatic lapse rates, respectively.  $\theta_e$  is determined using the Bolton  
191 (1980) approximation for the Clausius-Clapeyron relationship and the method of Bryan  
192 (2008), which contains assumptions that are particularly accurate in heavily precipitating  
193 situations.

194 Once all coefficients and forcing have been determined, successive over-relaxation (SOR)  
195 is used to converge on a solution for the Sawyer-Eliassen streamfunction. Since (1) is a

196 second-order elliptic differential equation, a unique solution is guaranteed only when the  
197 quasi-geostrophic potential vorticity (QGPV) is greater than zero at each grid point. There-  
198 fore, in order to facilitate convergence, if any grid point has QGPV less than zero, the SOR  
199 algorithm calculates a 4-point average of the Sawyer-Eliassen streamfunction at the neigh-  
200 boring grid points and assigns the averaged value to the grid point of interest during each  
201 iteration. For the solutions<sup>3</sup> presented here, the ageostrophic streamfunction is set to zero  
202 on the boundaries of the cross section, as in the solutions presented by Todsén (1964) and  
203 Shapiro (1981).

204     Employing (1), Shapiro (1982) demonstrated that, in the absence of along-jet geostrophic  
205 temperature advection, the ageostrophic circulations associated with geostrophic stretching  
206 deformation resembled the traditional four-quadrant model with a thermally direct (indirect)  
207 circulation in the jet entrance (exit) region (Fig. 5a). Along-jet geostrophic temperature  
208 advection mobilizes the geostrophic shearing deformation forcing, which ‘shifts’ the thermally  
209 direct (indirect) circulation to the anticyclonic (cyclonic) shear side of the jet for cases of  
210 geostrophic cold air advection, such that subsidence is present through the jet core (Fig.  
211 5b). Conversely, geostrophic warm air advection along the jet acts to ‘shift’ the thermally  
212 direct (indirect) circulation to the cyclonic (anticyclonic) shear side of the jet such that  
213 ascent occurs through the jet core (Fig. 5c)<sup>4</sup>. These vertical motions have been shown  
214 by numerous studies to exert a considerable influence on restructuring the tropopause and  
215 baroclinicity in the upper troposphere and lower stratosphere (e.g. Reed and Sanders 1953;  
216 Reed 1955; Shapiro 1981; Shapiro 1982; Keyser and Pecnick 1985; Lang and Martin 2012)  
217 and can affect the production of sensible weather.

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<sup>3</sup>Inspection of the computed ageostrophic circulation in cross sections immediately upstream and down-  
stream of those selected for the forthcoming analysis were extremely consistent.

<sup>4</sup>These circulations can also be understood in terms of positive/negative vorticity advection by the thermal  
wind (i.e. Sutcliffe 1947) as described by Martin (2014).

### 218 **3. The 1-2 May 2010 Nashville Flood - Overview**

219 M12 and Durkee et al. (2012) provide excellent overviews of the meso- and synoptic-scale  
220 processes responsible for the production of precipitation across the region and the reader is  
221 referred to those works for any additional details. Here, we present an abbreviated synoptic  
222 overview of the period from 0000 UTC 1 May - 0000 UTC 2 May across the contiguous  
223 United States.

224 Figure 6 illustrates the anomalous nature of the moisture that was in place during the  
225 week of the event and shows that most of the region east of the Mississippi River was  
226 characterized by precipitable water values that were at least 5 mm greater than normal for  
227 late April/early May. The figure also captures the filamentary structure of the anomaly  
228 pattern over the Gulf of Mexico, typical of an atmospheric river. Furthermore, Nashville  
229 observed a precipitable water value of 51.3 mm (2.02 in.) at 0000 UTC 2 May, registering  
230 well above the 99th percentile for that time of year (45.7 mm) and indicating an almost  
231 unprecedented availability of moisture in the troposphere throughout the duration of the  
232 flooding event (NWS 2011).

233 The large-scale pattern at 0000 UTC 1 May (Fig. 7a) depicted an occluding mid-latitude  
234 cyclone, with a sea-level pressure (SLP) minimum below 988 hPa, located along the North  
235 Dakota/Manitoba border. A warm front at the surface extended across the northern Great  
236 Lakes eastward towards New York, while a cold front stretched from northeastern Minnesota  
237 southward into eastern Texas. Immediately to the east of the cold front was a tongue of  
238 poleward moisture flux at 925 hPa, which flowed from the Gulf of Mexico into the Great  
239 Lakes. Maximum poleward moisture flux values<sup>5</sup> over the northern Gulf of Mexico were  
240 greater than  $30 \text{ cm s}^{-1}$  at this time along the axis of maximum moisture flux<sup>6</sup>. At 250 hPa,

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<sup>5</sup>Poleward moisture flux is computed as the product of the y-direction velocity,  $v$  ( $\text{m s}^{-1}$ ), and the mixing ratio ( $\text{kg kg}^{-1}$ ). Given typical values for mixing ratio ( $5 \text{ g kg}^{-1}$ ) and wind speed ( $10 \text{ m s}^{-1}$ ), this calculation would yield a moisture flux of  $.05 \text{ m s}^{-1}$  or  $5 \text{ cm s}^{-1}$ .

<sup>6</sup>Defined as the axis of maximum convergence of the moisture flux gradient vector, it is included as a common reference point for determining the impact of the forthcoming Sawyer-Eliassen circulations on the

241 a polar jet, as denoted by the blue arrow, stretched from Baja California northeastward into  
242 the Central Plains in association with a deep upper-level trough over the western United  
243 States, while a subtropical jet, as identified by the red arrow, extended across northern  
244 Mexico and eastward along the Gulf Coast. At this time, note that even though the two jets  
245 were in relatively close proximity to one another, they were not superposed.

246 By 0000 UTC 2 May (Fig. 7b), the mid-latitude cyclone had remained stationary along  
247 the Canadian border and had begun to decay. The cold front, while making slight progress  
248 to the east over the Great Lakes, was stationary over portions of the southern Mississippi  
249 River Valley, helping to focus precipitation over the same areas for a second consecutive day.  
250 Notably, the poleward moisture flux at 925 hPa was substantially larger than at the earlier  
251 time, with maximum values over the northern Gulf of Mexico now exceeding  $40 \text{ cm s}^{-1}$  along  
252 the axis of maximum moisture flux. Coincident with this increase in moisture flux was the  
253 first indication of a jet superposition in the upper troposphere over portions of southwestern  
254 Oklahoma and western Texas, as denoted by the purple line. This jet superposition was  
255 characterized by a rapid acceleration of the jet core wind speeds, which exceeded  $70 \text{ m s}^{-1}$   
256 at this time.

257 A cross section along the line C-C' in Fig. 7b, perpendicular to the jet core, is shown in  
258 Fig. 8 and confirms the presence of a superposed jet. Rather than the three-step tropopause  
259 structure described by DT57, this cross section is characterized by a two-step tropopause  
260 structure with a vertical PV wall that distinguishes the break between the polar (400 hPa)  
261 and the tropical tropopause (125 hPa). Also note that the identification criteria for the polar  
262 and the subtropical jet are both met within the same vertical grid column that intersects  
263 the jet core. The superposed jet is also associated with considerable upper tropospheric and  
264 lower stratospheric baroclinicity, as required to support the increased vertical wind shear.  
265 The coincidence of the observed increase in poleward moisture flux with a proximate jet  
266 superposition event suggests that the ageostrophic circulation associated with the superposed  

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poleward moisture flux.

267 jet may have played a role in the increased poleward moisture flux observed over the southern  
268 Mississippi River Valley.

269 To investigate this possibility, Fig 9a depicts the total change in the 925 hPa poleward  
270 moisture flux across the southern Mississippi River Valley during the 24-h period from 0000  
271 UTC 1 May - 0000 UTC 2 May. Results demonstrate that the poleward moisture flux  
272 increases by roughly  $9 \text{ cm s}^{-1}$  south of the Gulf Coast in the vicinity of the axis of maximum  
273 moisture flux at both times. This is in general agreement with the qualitative assessment  
274 made from Fig. 7. However, an examination of the difference in the geostrophic poleward  
275 moisture flux over the same period (Fig. 9b) shows little to no change along the axes of  
276 maximum moisture flux. Instead, increased geostrophic fluxes are displaced to the north and  
277 east, consistent with a shift of the strongest southerly geostrophic winds in that direction.  
278 So, while M12 and Durkee et al. (2012) note that the largest fraction of the moisture flux  
279 was accomplished by geostrophic processes during this event, the majority of the observed  
280 *increase* in moisture flux south of New Orleans is accounted for by changes in the ageostrophic  
281 poleward moisture flux (which includes the effects of the jet circulation, as well as curvature,  
282 friction, etc). Figure 9c confirms this notion, depicting an increase on the order of  $9 \text{ cm}$   
283  $\text{s}^{-1}$  along the Gulf Coast and centered squarely on the axes of maximum moisture flux.  
284 Given this conclusion, the analysis that follows aims to determine the specific impact of  
285 the superposed jet circulation on the ageostrophic moisture flux over the northern Gulf of  
286 Mexico.

## 287 4. Diagnosis of Sawyer-Eliassen Circulations

288 The analysis begins with an investigation of the role the ageostrophic circulation associ-  
289 ated with the superposed jet played in facilitating poleward moisture flux into the southern  
290 Mississippi River Valley. The individual forcing terms for the superposed jet circulation are  
291 then examined to better understand their impacts on the resultant circulation.

293 At 0000 UTC 1 May, an area of convection was beginning to form over portions of central  
294 Arkansas. These thunderstorms would later move off to the east and form the first MCS  
295 that dropped considerable rainfall amounts across portions of the Tennessee River Valley  
296 on the first day of the event. Additionally, the polar and subtropical jets bifurcated over  
297 northern Texas, with the polar jet extending to the northeast over the Central Plains while  
298 the subtropical jet stretched eastward along the Gulf Coast (Fig. 7a). As such, a diagnosis  
299 at this time must consider the separate ageostrophic circulations associated with each jet  
300 and its overall contribution to the poleward moisture flux across the southern Mississippi  
301 River Valley.

302 Figure 10a shows the Sawyer-Eliassen circulation along the cross section from D-D' in Fig.  
303 7a, which is cut through the subtropical jet's exit region and is nearly parallel to the axis of  
304 maximum moisture flux near the Gulf Coast. The solution depicts a rather weak thermally  
305 indirect circulation with the strongest upward vertical motions and streamfunction maximum  
306 centered close to Little Rock, Arkansas (LZK), largely associated with the diabatic effects of  
307 the ongoing convection (not shown). Over the northern Gulf of Mexico, where the poleward  
308 moisture flux was maximized at this time, the role of the Sawyer-Eliassen circulation is  
309 rather unimpressive, with a maximum contribution on the order of  $5 \text{ cm s}^{-1}$  around 925  
310 hPa. A comparison with the total observed ageostrophic poleward moisture flux at this  
311 time (Fig. 10b) shows that the magnitude of the poleward moisture flux associated with  
312 the Sawyer-Eliassen circulation is on par with observed ageostrophic flux values over the  
313 northern Gulf of Mexico. As a result, it is reasonable to conclude that our calculation  
314 accurately captures the maximum contribution to the overall moisture flux made by the  
315 subtropical jet's ageostrophic circulation in that region.

316 As previously indicated, the polar jet was located further to the north and west over  
317 the Central Plains. Figure 11 demonstrates that the Sawyer-Eliassen circulation associated  
318 with the polar jet, along the cross-section E-E' in Fig. 7a, is a stronger, thermally *direct*

319 circulation, such that the low-level, horizontal branch of this circulation actually opposes  
320 the poleward moisture flux promoted by the subtropical jet. However, the juxtaposition of  
321 these two circulations is favorable for promoting upward vertical motions directly over Little  
322 Rock, where the ascending branches of both circulations are collocated. Therefore, while  
323 the separate jet circulations likely played a symbiotic role in aiding the initial formation of  
324 convection that occurred over central Arkansas at 0000 UTC 1 May, the subtropical jet's  
325 circulation was the only one capable of facilitating a poleward moisture flux into the southern  
326 Mississippi River Valley at that time.

327 At 0000 UTC 2 May, an area of convection was ongoing over portions of southern  
328 Arkansas and northern Louisiana. As mentioned previously (and illustrated in Fig. 7),  
329 the poleward moisture flux increased considerably over the intervening 24-h to a maximum  
330 value greater than  $40 \text{ cm s}^{-1}$ , coincident with the jet superposition event. It is important  
331 to note that mixing ratios across the southern Mississippi River Valley and northern Gulf of  
332 Mexico were largely unchanged between the two days (not shown). As a result, the increase  
333 in poleward moisture flux was a direct consequence of an increase in wind speed. To investi-  
334 gate the impact of the superposed jet on the poleward moisture flux, we return to the cross  
335 section labeled C-C' in Fig. 7b, drawn perpendicular to the superposed jet axis and through  
336 the axis of maximum poleward moisture flux at 0000 UTC 2 May. The solution for the  
337 circulation within this cross section, shown in Fig. 12a, depicts a robust thermally indirect  
338 circulation, much stronger than the circulation associated solely with the subtropical jet at  
339 the previous time (Fig. 10a), and shifted towards the anticyclonic shear side of the jet. The  
340 superposed jet circulation is characterized by 1) a plume of ascent that extends from the  
341 surface through the jet core with local maxima found in both the middle and lower tropo-  
342 sphere, and 2) much stronger moisture fluxes over the northern Gulf of Mexico, maximized  
343 around  $15 \text{ cm s}^{-1}$  near 925 hPa.

344 The cross section C-C' is oriented at an angle to the axis of maximum moisture flux at this  
345 time. Consequently, in order to facilitate a direct comparison between the poleward moisture

346 fluxes associated with both the subtropical and the superposed jets, the component of the  
347 moisture flux associated with the superposed jet in the direction of the axis of maximum  
348 moisture flux at 0000 UTC 2 May was calculated and determined to be  $11 \text{ cm s}^{-1}$  at 925  
349 hPa. This is an increase of about  $6 \text{ cm s}^{-1}$  (a 120% increase) from that associated solely  
350 with the subtropical jet at the earlier time. Figure 12b shows that this value is on par  
351 with, but slightly larger than, observed ageostrophic poleward moisture fluxes just south of  
352 New Orleans. This overestimate is at least partially a result of the fact that we neglect the  
353 effects of friction and flow curvature on the ageostrophic circulation in our solution of the  
354 Sawyer-Eliassen equation. Recalling that total ageostrophic moisture flux values increased  
355 by as much as  $9 \text{ cm s}^{-1}$  over the 24-h period (Fig. 9c), we conclude that the ageostrophic  
356 circulation associated with superposed jet accounts for the vast majority of the increased  
357 poleward moisture flux. As demonstrated by M12, this moisture flux was crucial in the  
358 production of precipitation further to the north during the flooding event. Thus, the analysis  
359 presented here illustrates the role the intensified Sawyer-Eliassen circulation associated with  
360 the superposed jet played in magnifying the severity of the event.

361 *b. Partition of Sawyer-Eliassen Forcing Terms*

362 The diagnostic power of the Sawyer-Eliassen equation (1) lies in the fact that the forcing  
363 can be broken down into the separate geostrophic forcing terms (shearing and stretching  
364 deformation) and a diabatic term. Consequently, the circulation associated with the super-  
365 posed jet can be further dissected in order to gauge the significance of the respective forcing  
366 terms in shaping its sense and strength. The portion of the Sawyer-Eliassen circulation asso-  
367 ciated with the total geostrophic forcing ( $Q_g$ ) is shown in Fig. 13a and depicts a circulation  
368 that, similar to the full circulation (Fig. 12a), is thermally indirect and shifted towards the  
369 anticyclonic shear side of the jet, positioning ascent directly beneath the jet core.

370 Intriguing differences, however, are found when comparing the distribution of the vertical  
371 motion and moisture flux to that shown in Fig. 12a. In contrast to the full circulation (Fig.

372 12a), which has a plume of ascent from the surface through the jet core, the  $Q_g$  circulation  
373 (Fig. 13a) has its strongest vertical motions primarily confined to the middle and upper  
374 troposphere. In addition, the low-level, horizontal branch of the  $Q_g$  circulation near the  
375 surface at the Gulf Coast is far weaker, with low-level moisture flux values only around 3  
376  $\text{cm s}^{-1}$ , much smaller than those forced by the full circulation.

377 The  $Q_g$  circulation, in Fig. 13a, can be partitioned into the individual circulations  
378 associated with the geostrophic shearing ( $Q_{SH}$ ) and stretching ( $Q_{ST}$ ) deformation terms,  
379 respectively. The  $Q_{SH}$  circulation is shown in Fig. 13b and depicts a thermally indirect  
380 circulation that is positioned primarily on the cyclonic shear side of the jet. This places the  
381 descending branch of the circulation directly beneath the jet core, opposite to the ascent  
382 observed in that region in the  $Q_g$  circulation. Examination of both the temperature gradient  
383 and geostrophic wind normal to the cross section suggests that areas between roughly 400-  
384 800 hPa were characterized by geostrophic cold air advection in cyclonic shear ( $Q_{SH} < 0$ ),  
385 consistent with the thermally indirect characteristics of the circulation observed in Fig. 13b.

386 Figure 13c shows that  $Q_{ST}$  acts to drive a thermally direct circulation about the strong  
387 upper tropospheric front centered on the cyclonic shear side of the jet, but offset slightly  
388 poleward of the center of the  $Q_{SH}$  circulation (Fig. 13b). Consequently,  $Q_{ST}$  promotes ascent  
389 directly beneath the jet core, slightly poleward of, and thus counteracting, the subsidence  
390 associated with  $Q_{SH}$ . Investigation of the along-cross section geostrophic wind shows a region  
391 of geostrophic confluence centered squarely on the upper tropospheric front ( $Q_{ST} > 0$ ),  
392 which would act to enhance the horizontal temperature gradient around 500 hPa and drive  
393 a thermally direct circulation.

394 Interestingly, this cross section is drawn through a geostrophic jet exit region at 500 hPa,  
395 as shown in Fig. 14. Typically, such regions are characterized by diffluent flow and associated  
396 horizontal frontolysis in the vicinity of any regions of baroclinicity, resulting in a thermally  
397 indirect circulation. Figure 14 shows that in this case, an embedded shortwave trough over  
398 the panhandles of Oklahoma and Texas actually produces a region of geostrophic confluence

399 in the vicinity of the geostrophic jet exit region. This confluence is responsible for an area  
400 of horizontal geostrophic frontogenesis precisely in the location in which a thermally direct  
401 circulation is observed in Fig. 13c.

402 Comparison of the intensities and areal extents of the  $Q_{SH}$  and  $Q_{ST}$  circulations demon-  
403 strates that the  $Q_{SH}$  circulation is the dominant component. Consequently, the sum of the  
404 two circulations indicates that the  $Q_{ST}$  circulation acts to erode the updraft associated with  
405 the  $Q_{SH}$  circulation on the cyclonic shear side of the jet, while preserving the downdraft on  
406 the anticyclonic shear side. The net result remains a thermally indirect circulation, but one  
407 that is shifted towards the anticyclonic shear side of the jet with ascent directly beneath the  
408 jet core. This total  $Q_g$  circulation is displaced further equatorward than might be expected  
409 under a regime of geostrophic cold air advection in cyclonic shear within a geostrophic jet  
410 exit region (Fig. 5b) due to the effects of the geostrophic confluence associated with the  
411 shortwave trough.

412 The final contribution to the full Sawyer-Eliassen circulation comes from the diabatic  
413 forcing. Figure 15a shows that the circulation associated with the diabatic forcing is focused  
414 entirely below 400 hPa, where latent heating acts to produce a dipole centered slightly  
415 north of the Gulf Coast, with a thermally direct circulation further to the north and a  
416 stronger thermally indirect circulation to the south. Upward vertical motions associated with  
417 this diabatically-induced circulation are also focused in the lower troposphere and coincide  
418 well with the area of most intense latent heat release from the initial convective activity.  
419 Most notably, the poleward moisture flux associated with the thermally indirect diabatic  
420 circulation (Fig. 15a) is much stronger than that associated with the  $Q_g$  circulation (Fig.  
421 13a), with values greater than  $9 \text{ cm s}^{-1}$  over the northern Gulf of Mexico. Consequently,  
422 the majority of the poleward moisture flux produced by the full ageostrophic circulation is  
423 driven by the diabatic component.

424 The preceding discussion indicates that the  $Q_g$  forcing largely determines the mid-  
425 tropospheric portion of the full Sawyer-Eliassen circulation (Fig. 12a). The diabatic portion,

426 then, provides a means by which the full tropospheric-deep circulation communicates directly  
427 with the surface, as it was responsible for the majority of the increase in low-level poleward  
428 moisture flux into the southeast United States and also coupled surface-based vertical mo-  
429 tions to those in the middle troposphere.

430 The analysis also suggests a crucial positive feedback mechanism that, on its own, may  
431 act to further strengthen and promote the longevity of the entire Sawyer-Eliassen circulation.  
432 Strong moisture flux and subsequent ascent promotes latent heat release through condensa-  
433 tion. The latent heat release produces a lower tropospheric ageostrophic circulation that can  
434 further strengthen the poleward moisture flux into a region and, subsequently, increase the  
435 potential for additional latent heat release. The addition of middle and upper tropospheric  
436 ascent provided by the  $Q_g$  circulation to that induced by the diabatically-forced circulation  
437 promotes the vigorous and tropospheric-deep vertical motions necessary for the production  
438 of heavy precipitation and intense latent heat release. In addition, the strong latent heat  
439 release beneath the jet core can act to erode upper-level PV, helping to fortify the vertical  
440 PV wall associated with the superposed jet structure thereby acting to maintain, or even  
441 strengthen, the strong wind speeds that are associated with it.

442 Support for the veracity of the superposed jet's diagnostic Sawyer-Eliassen circulation  
443 is also evident in the cross section of vertical motion from the GFS analysis, shown in Fig.  
444 15b. Similar to the tropospheric-deep plume of ascent observed with the full superposed  
445 jet circulation in Fig. 12a, the GFS shows a continuous plume of ascent that runs roughly  
446 parallel to the leading edge of the upper tropospheric front and through the jet core. In  
447 addition, the distribution of vertical motion depicts two local maxima, one near the Gulf  
448 Coast in the vicinity of the maximum latent heat release (Fig. 15a), and another in the  
449 middle-to-upper troposphere that is nearly collocated with the maximum in ascent associated  
450 with the  $Q_g$  portion of the superposed jet circulation (Fig. 13a).

451 A similar positive feedback mechanism, envisioned from a PV perspective, was proposed  
452 by Lackmann (2002) in his study of a warm conveyor belt during a February 1997 cyclogenesis

453 event and serves as an analog to the mechanism discussed above. In that case, it was found  
454 that the circulation associated with a linear, diabatically-generated positive PV anomaly  
455 along a low-level frontal boundary made a non-negligible contribution to the strength of the  
456 southerly low-level jet. The strengthened low-level jet then accomplished additional poleward  
457 moisture transport into the region, further conditioning the atmosphere for additional latent  
458 heat release. Indeed, Lackmann (2013) found similar conditions at work during the Nashville  
459 flood, where the low-level jet was characterized by a linear positive PV anomaly to its west,  
460 along the stationary cold frontal boundary. While the study indicated that topographic  
461 effects along the Mexican Plateau were the primary mechanism behind the initial generation  
462 of low-level cyclonic PV present along the frontal boundary during the event, diabatic effects  
463 acted to enhance the magnitude of these anomalies as they drifted eastward into the southern  
464 United States.

## 465 **5. Discussion and Conclusions**

466 The analysis presented here demonstrates that the lower tropospheric horizontal branch  
467 of the Sawyer-Eliassen circulation associated with a superposed jet helped to enhance the  
468 poleward moisture flux prior to the second day of the 2010 Nashville flood event. This  
469 explanation accounts for the analyses by M12 and Durkee et al. (2012) and their particular  
470 observations of increased poleward moisture transport during the second day of the event.  
471 Mixing ratios on these two days were largely unchanging across the southern Mississippi River  
472 Valley. Given this fact, an increased wind speed underlies the increased poleward moisture  
473 flux that was observed on the second day. The analysis presented here shows that this  
474 increased wind speed is primarily attributable to the ageostrophic circulation associated with  
475 the superposed jet and illuminates one mechanism by which such superposed jet structures  
476 may have an influence on the evolution of a high-impact weather event. Such a dynamical  
477 influence is undoubtedly magnified by the fact that the superposed circulation, by virtue of

478 its association with the subtropical jet, is able to draw upon the moist and weakly stratified  
479 air mass characteristic of the lower troposphere equatorward of the subtropical jet.

480 Additionally, partition of the forcings driving the superposed jet circulation provides  
481 insights into its internal dynamics. In the case presented here, the  $Q_{SH}$  term was more dom-  
482 inant than the  $Q_{ST}$  term. As a result, the entire  $Q_g$  circulation took on the thermally indirect  
483 characteristics of the  $Q_{SH}$  circulation. The thermally direct circulation associated with the  
484  $Q_{ST}$  forcing, however, acted to significantly counteract the  $Q_{SH}$  circulation on the cyclonic  
485 shear side of the jet, shifting the locus of the entire  $Q_g$  circulation towards the anticyclonic  
486 shear side of the jet. Such an orientation can dynamically assist convection, as upward verti-  
487 cal motions on the anticyclonic shear side of the jet are exhausted in an area with much lower  
488 inertial stability. In comparison to the cases examined by Shapiro (1981) and Shapiro (1982),  
489 this observed circulation is atypical for an environment of geostrophic cold air advection in  
490 a geostrophic jet exit region. It is important to note, however, that throughout much of the  
491 evolution of a superposed jet structure, the environment is characterized by more than one  
492 jet core. Therefore, idealized models of transverse circulations in environments characterized  
493 by single jet cores may not be expected to represent the circulations characterizing the more  
494 complex superposed jet environment.

495 Given that superposed jets are often characterized by anomalously strong wind speeds in  
496 the jet core, it is likely that the horizontal shear is also anomalously large in the vicinity of  
497 these features. Consequently, it is conceivable that the  $Q_{SH}$  term may consistently dominate  
498 the  $Q_g$  forcing for ageostrophic circulations associated with superposed jets, particularly  
499 away from geostrophic jet entrance and exit regions. A more comprehensive examination  
500 of other superposed jet streaks may illuminate the nature of the interaction between the  
501 two geostrophic forcing terms in the vicinity of these structures and how their circulations  
502 compare with established conceptual models.

503 Moreover, this case illustrates that latent heat release can have a considerable impact on  
504 shaping and enhancing the entire ageostrophic circulation. If they couple favorably, the  $Q_g$

505 and diabatic circulations can drive a notable positive feedback mechanism, similar to that  
506 proposed by Lackmann (2002), which can act to both strengthen upward vertical motions  
507 and intensify the ageostrophic winds in the low-level horizontal branch of the circulation.  
508 Studies of jet circulations in other heavy precipitation events may help to further characterize  
509 this feedback mechanism.

510 Both  $Q_g$  circulations and latent heat release also have the ability to reshape the tropopause  
511 and, subsequently, affect the structure of the jet. The development of superposition events,  
512 in particular, is usually characterized by 1) the melding of two separate tropopause folds as-  
513 sociated with the separate jets into a single, steeper one and 2) an attendant acceleration of  
514 the jet. We hypothesize that these transformations result from an interaction between latent  
515 heat release and internal jet-front dynamics. Figure 16 shows an idealized schematic in which  
516 both the polar and the subtropical jets are characterized by separate positive tropopause  
517 PV anomalies and cyclonic circulations that, while maximized at the level of the respective  
518 anomalies, extend vertically through their respective columns. When these anomalies come  
519 in close proximity to one another, their circulations can interact with a potential for de-  
520 structive interference in the space between the separate jet cores, diminishing wind speeds  
521 in that location. Constructive interference in the column previously located between the two  
522 jets occurs when the two anomalies become vertically superposed and the tropopause steep-  
523 ens, producing the rapid acceleration of jet wind speeds and intensification of the secondary  
524 circulation that may be characteristic of a superposition event.

525 Latent heating from lower tropospheric frontal convection, which erodes upper-level PV,  
526 can promote steepening of the tropopause on the equatorward side of the subtropical jet. At  
527 the same time, stratospheric geostrophic warm air advection in cyclonic shear, as was present  
528 in the Nashville case both before and at the time of jet superposition (Fig. 13b), promotes  
529 ascent through the jet core via the  $Q_{SH}$  term, with subsidence poleward of the subtropical  
530 jet's exit region in the lower stratosphere (Lang and Martin 2012). This subsidence can act to  
531 flatten the PV trough in the space between the two jets, resulting in a single, more intensely

532 sloped tropopause, the superposition of the two PV anomalies (jets), and subsequent, rapid  
533 accelerations in jet wind speeds and the attendant ageostrophic circulations. The initial  
534 latitudinal separation of the two jet cores is hypothesized to be critical in this process. If  
535 they are so far apart as to place the lower tropospheric convection between the two jets,  
536 the convection may actually inhibit superposition by strengthening and reinforcing the PV  
537 trough between the two jets. Work is ongoing to investigate the nature of the processes that  
538 contribute to jet superposition and to what degree these processes are dependent upon, and  
539 sensitive to, characteristic distributions of latent heat release.

540 Consideration of this problem via the piecewise PV inversion scheme of Davis and  
541 Emanuel (1991) is also currently underway. We are devising analysis schemes by which  
542 the respective PV anomalies associated with both the polar and subtropical jets can be iso-  
543 lated and individually inverted in order to determine the circulations associated with each  
544 anomaly. Interactions between these separate circulations, and their individual ability to  
545 reshape the tropopause into the two-step structure characteristic of jet superpositions, will  
546 provide considerable insight into the process of superposition. The role of latent heat and  
547 surface-based PV anomalies can, in a similar manner, be interrogated from the PV perspec-  
548 tive, assisting in the development of a comprehensive picture of the dynamics driving jet  
549 superpositions. The results from this particular case study demonstrate that such features  
550 can, indeed, play a central role in the evolution of high-impact weather events. Consequently,  
551 greater understanding of the processes that conspire to form superposed jet structures, via  
552 consideration of internal jet-front dynamics from either the basic-state variables or PV per-  
553 spectives, can better inform forecasters regarding both the operation of such features as well  
554 as anticipation of their impacts.

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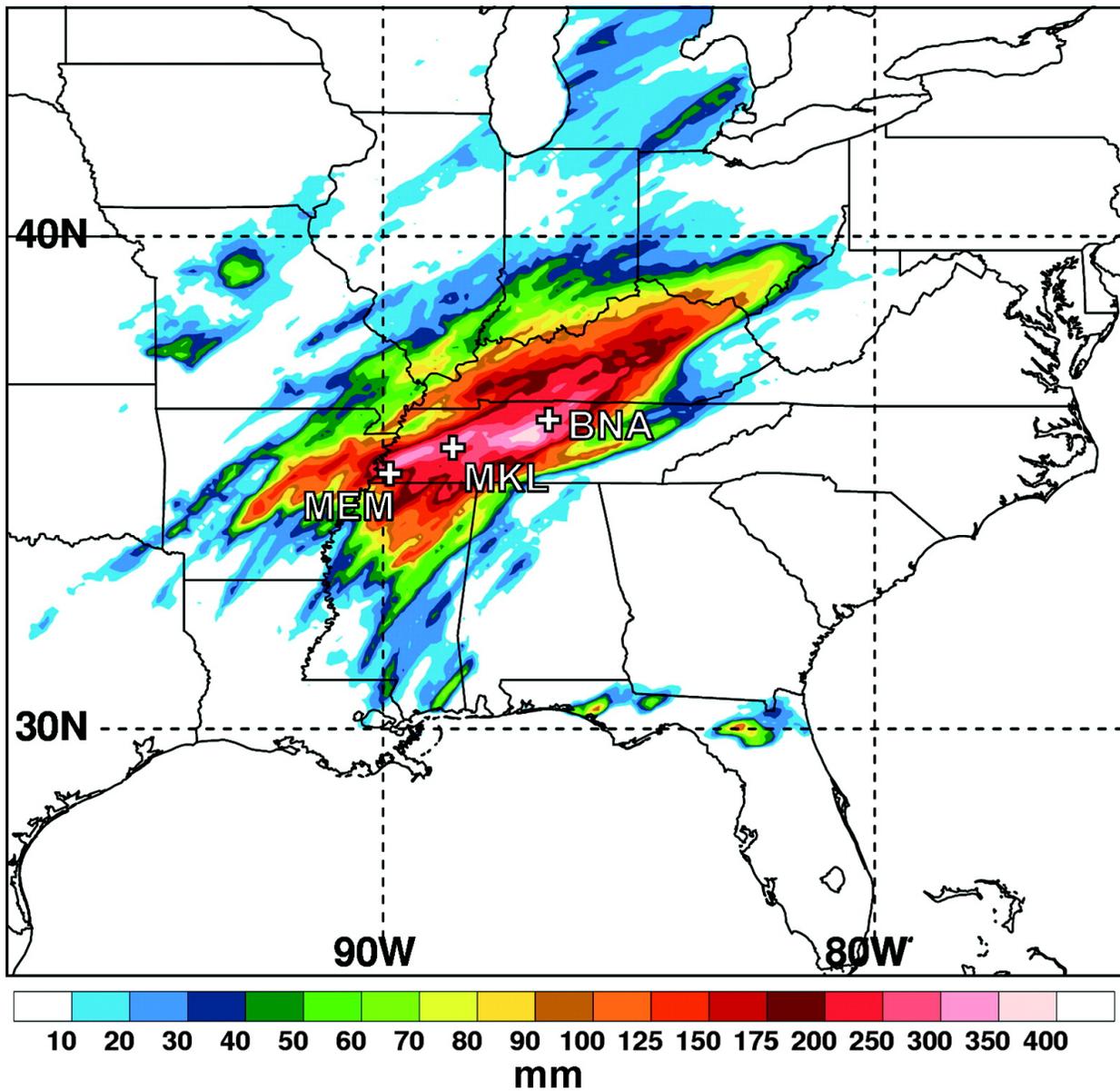


FIG. 1. 48-h precipitation estimates (shaded in mm following the color bar) for 0000 UTC 1 May - 0000 UTC 3 May 2010 from the National Precipitation Verification Unit quantitative precipitation estimates product. The location of Nashville (BNA), as discussed in the text, is identified. [From Moore et al. (2012; their Fig. 2)]

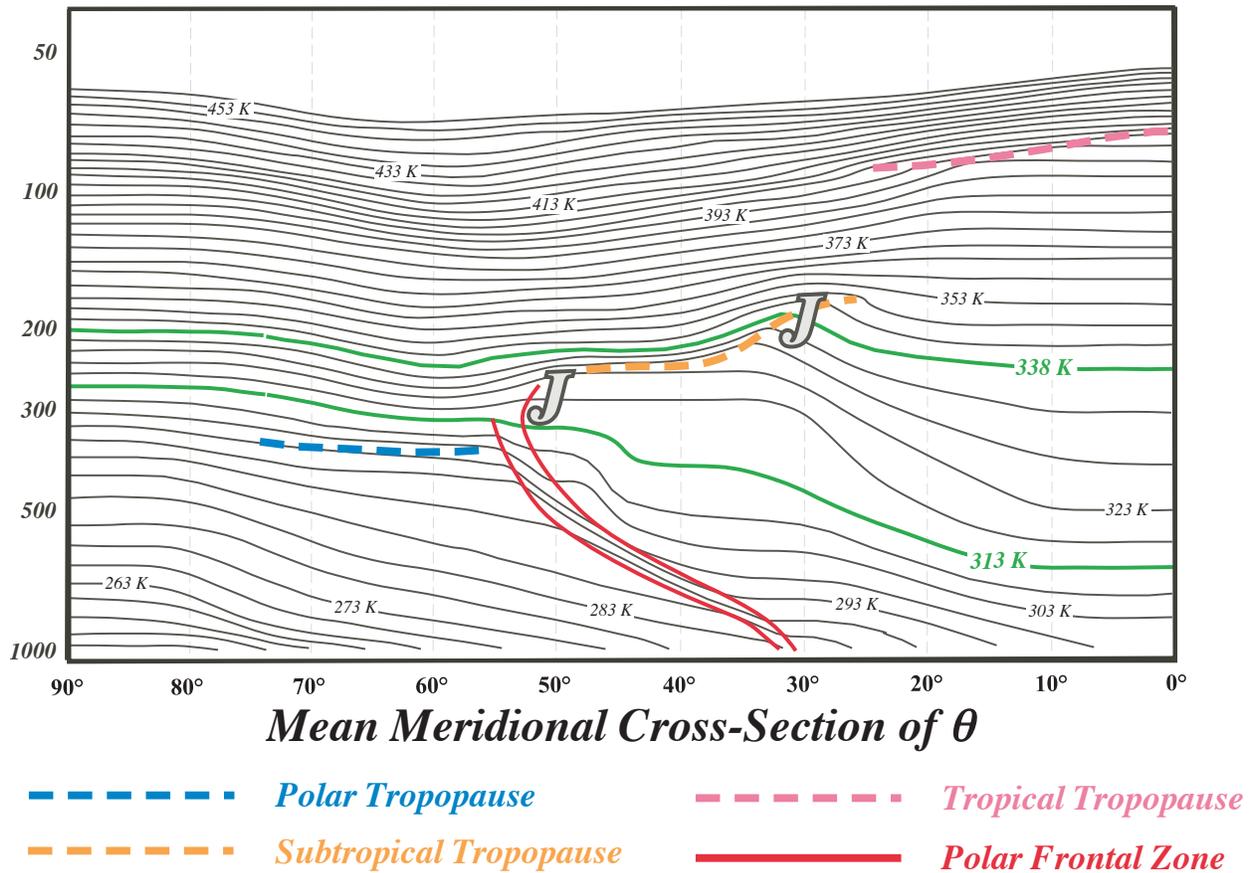


FIG. 2. Mean meridional cross section of potential temperature for 1 January 1956 with the polar, subtropical, and tropical tropopauses and polar frontal zone labeled as indicated in the legend. (Modified from Defant and Taba 1957; Fig. 13)

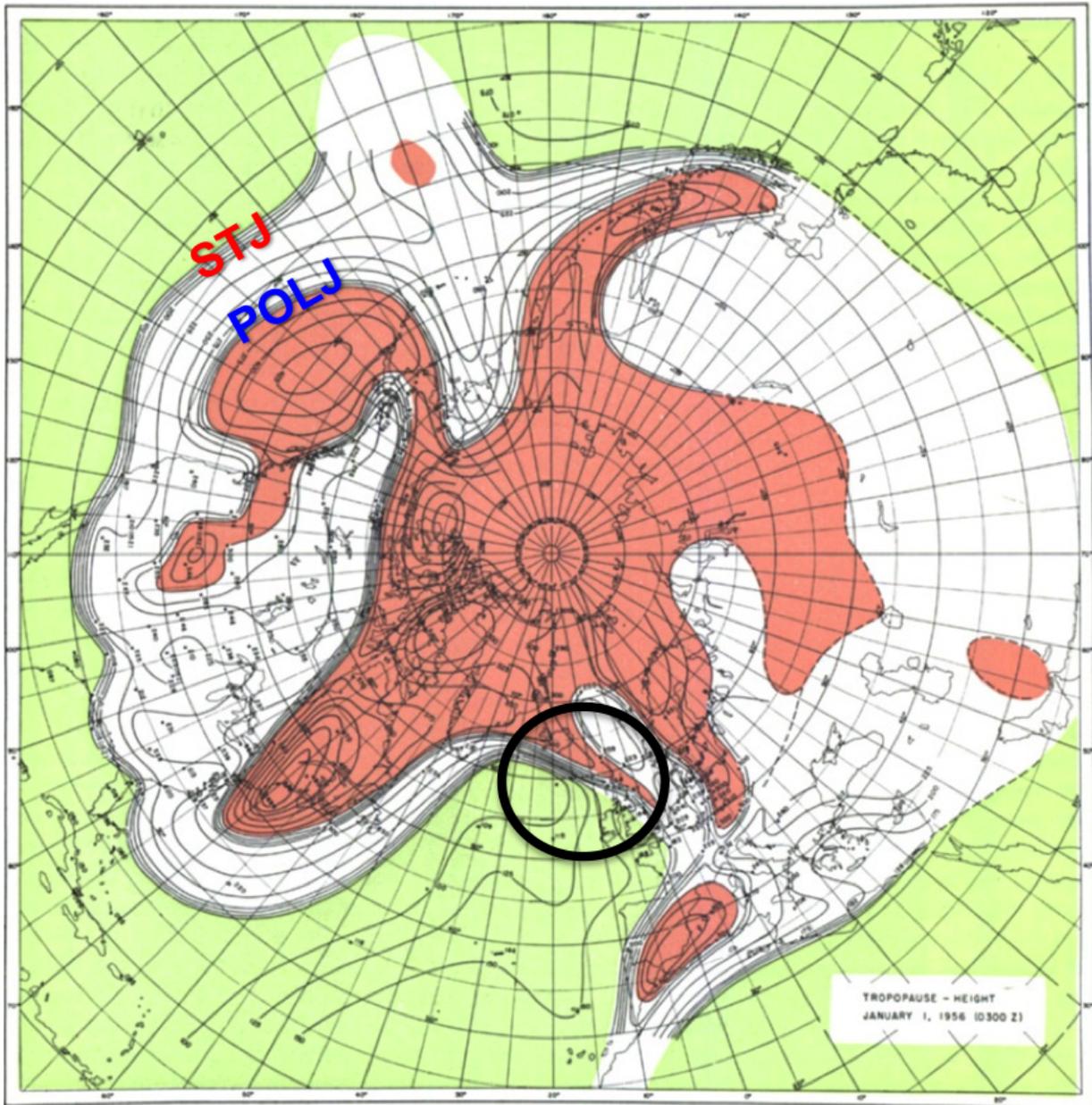


FIG. 3. Northern hemispheric map of tropopause height (hPa) at 0300 UTC 1 January 1956. Breaklines are denoted as areas with a sharp gradient in tropopause height. Breaklines that correspond to the subtropical (STJ) and polar jet (POLJ) are labeled accordingly. Area denoted with a circle is a region characterized by a superposition of the polar and subtropical jet. Green shading corresponds to the tropical tropopause, white shading the subtropical tropopause, and red the polar tropopause. (Modified from Defant and Taba 1957; Fig. 2)

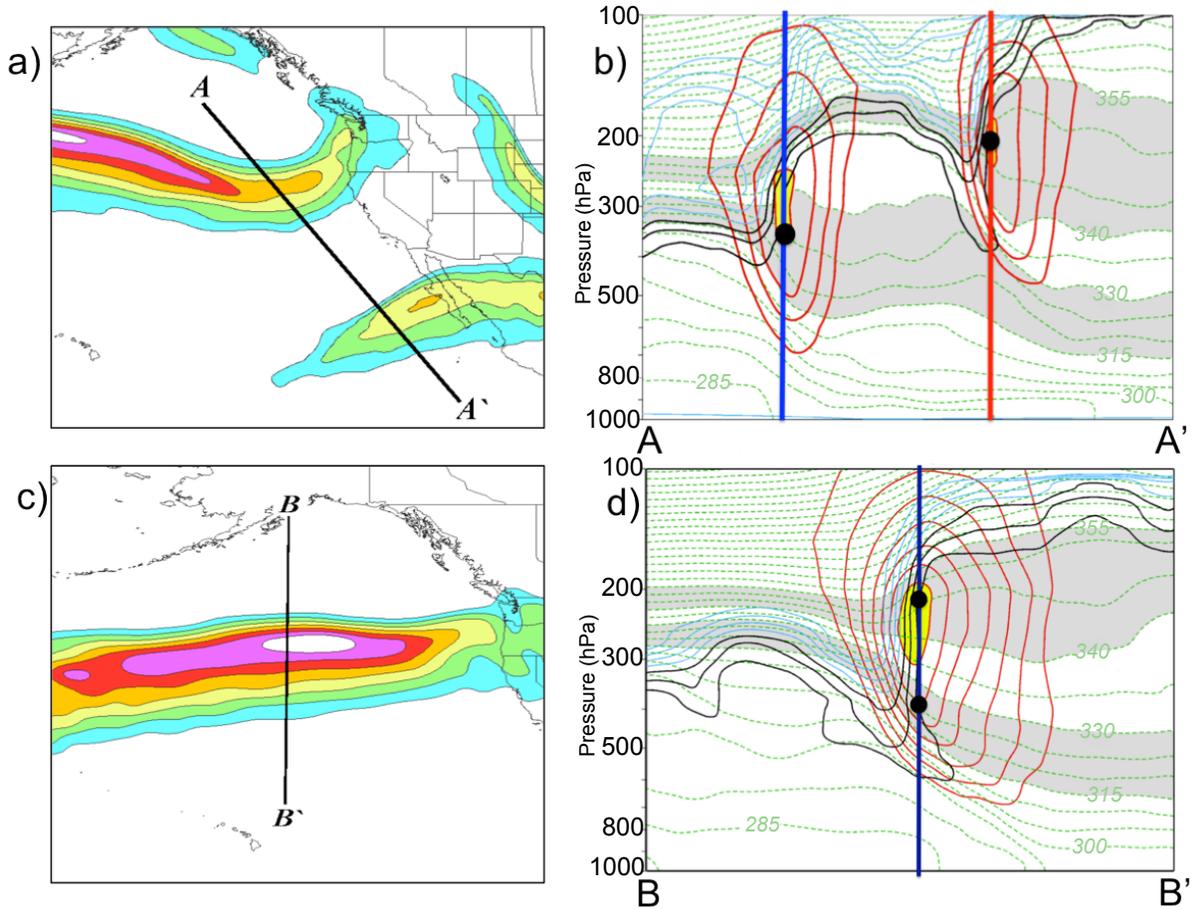


FIG. 4. (a) 300 hPa isotachs (shaded every  $10 \text{ m s}^{-1}$  starting at  $30 \text{ m s}^{-1}$ ) at 0000 UTC 27 April 2010 depicting separate polar and subtropical jets. (b) Cross section from A-A', in Fig. 4a, through separate polar and subtropical jet cores with contours of the 1,2,3 PVU ( $1 \text{ PVU} = \text{K m}^2 \text{ kg}^{-1} \text{ s}^{-1}$ ) surfaces (black), 4,5,6,7,8,9, PVU surfaces (light blue), potential temperature every 5K (dashed green), and isotachs every  $10 \text{ m s}^{-1}$  beginning at  $30 \text{ m s}^{-1}$  (red). The jet core is shaded yellow and the 315-330K and 340-355K isentropic layers, used to identify the location of the jets, are shaded gray. The blue (red) column corresponds to a grid column with the black dot confirming a positive ID of a polar (subtropical) jet. (c) Same as (a) depicting a superposed jet at 0000 UTC 24 October 2010. (d) Same as (b) but for the cross section from B-B', in Fig. 4c, with two positive IDs (black dots) within a single grid column indicating a jet superposition.

*Upper Troposphere*

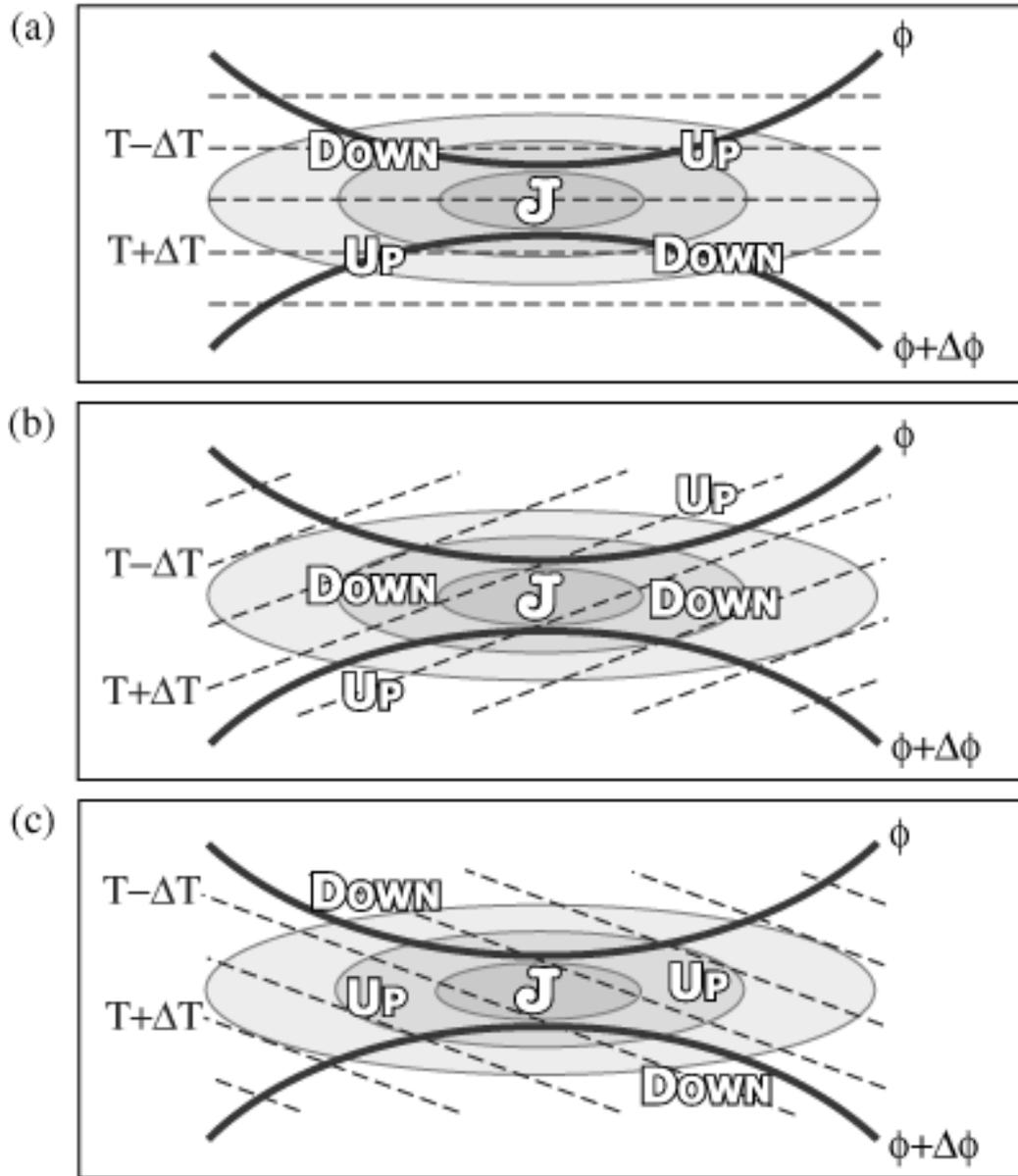


FIG. 5. Idealized configurations of jet circulations associated with a straight jet streak on an isobaric surface in the upper troposphere. Geopotential height (thick solid lines), potential temperature (dashed lines), geostrophic isotachs (fill pattern; with the jet speed maximum represented by the “J”), and Sawyer-Eliassen vertical motions indicated by “up” and “down” for a regime of (a) no geostrophic temperature advection (b) upper tropospheric geostrophic cold air advection and (c) upper tropospheric geostrophic warm air advection along the jet axis. (From Lang and Martin 2012; Fig. 3)

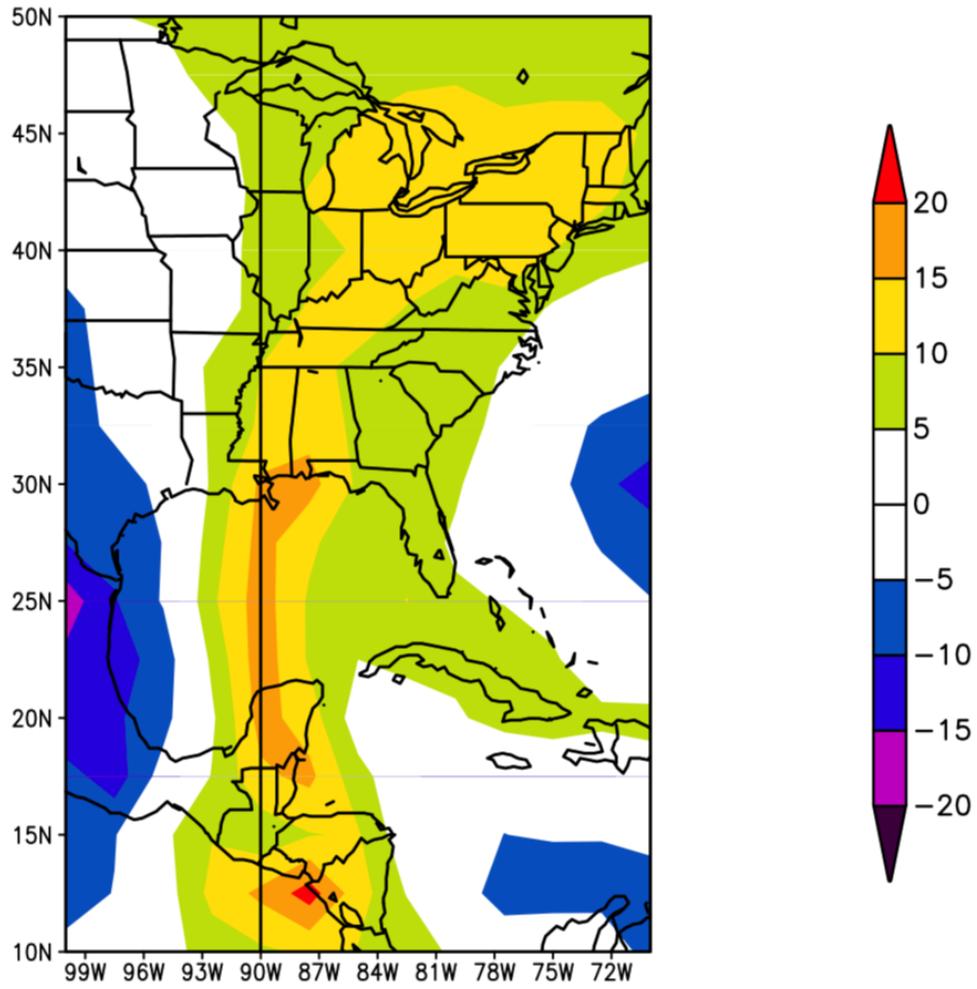


FIG. 6. 4-day precipitable water anomalies in mm (fill pattern) during the period of 30 April 2010 - 3 May 2010 across the eastern United States. (Earth Systems Research Lab)

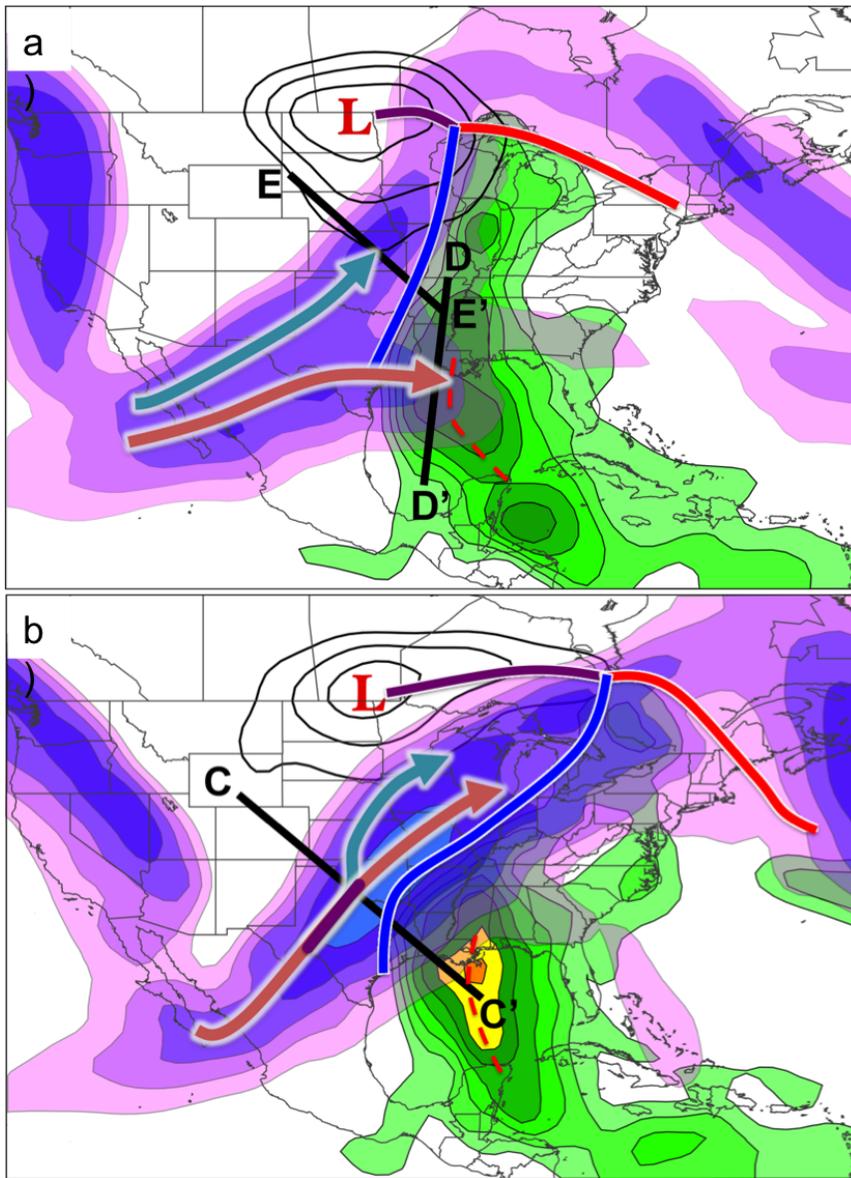


FIG. 7. Synoptic overview with sea-level pressure every 4 hPa beginning at 996 hPa (thin black lines), surface low pressure center (red “L”), surface frontal boundaries with the cold front denoted by the blue line, the warm front in red, and occluded front in purple, magnitude of the 925 hPa poleward moisture flux every  $5 \text{ cm s}^{-1}$  beginning at  $10 \text{ cm s}^{-1}$  (green fill pattern), 250 hPa isotachs every  $10 \text{ m s}^{-1}$  beginning at  $30 \text{ m s}^{-1}$  (purple fill pattern), the location of the polar (blue arrow), subtropical (red arrow), and superposed (purple line) jet, as identified using the algorithm defined in the text, and the axis of maximum 925 hPa poleward moisture flux (red dashed line) at (a) 0000 UTC 1 May 2010 and (b) 0000 UTC 2 May 2010.

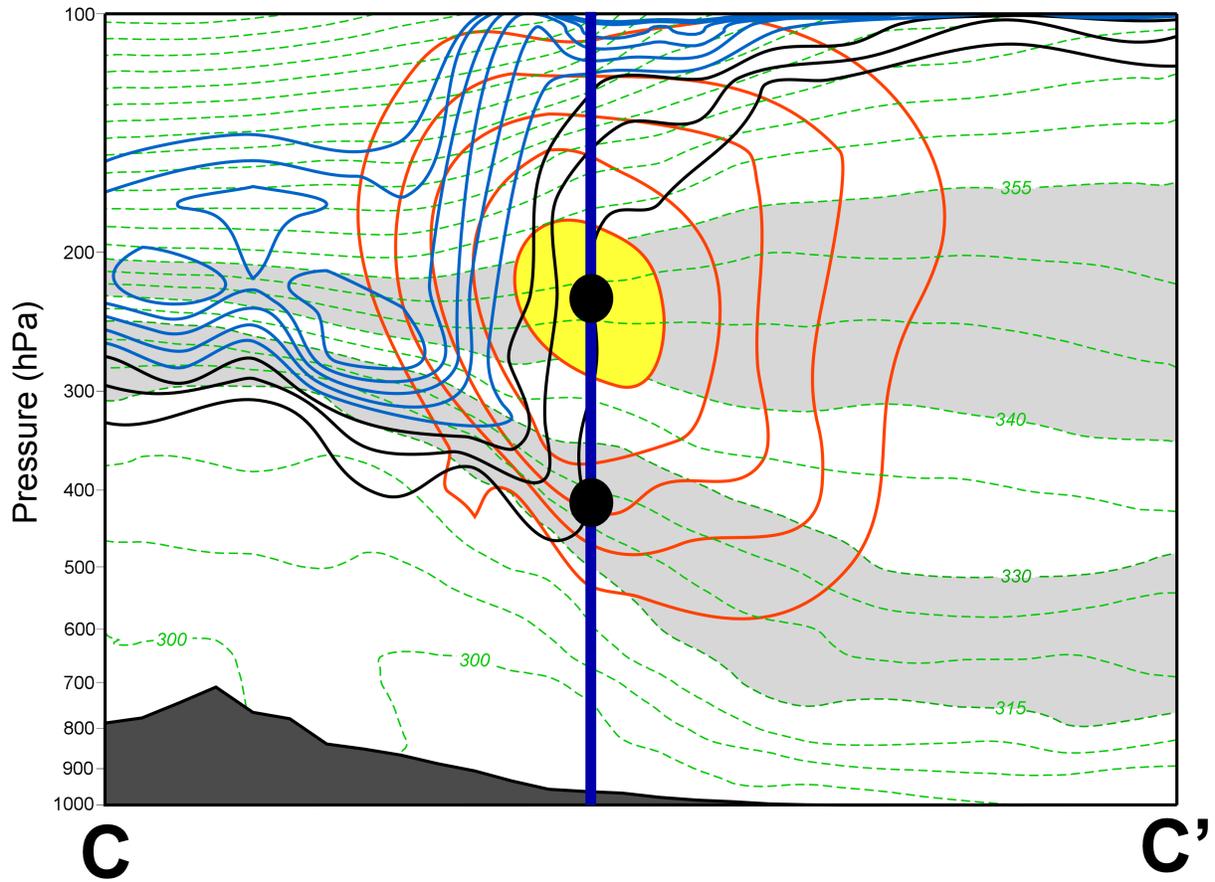


FIG. 8. Vertical cross section of potential temperature, potential vorticity, and isotachs at 0000 UTC 2 May 2010 along the line C-C' in Fig. 7b. Variables labeled, contoured, and shaded as in Fig. 4d. Black dots represent separate polar and subtropical jet identifications in the same grid column, which is identified by the bold vertical line.

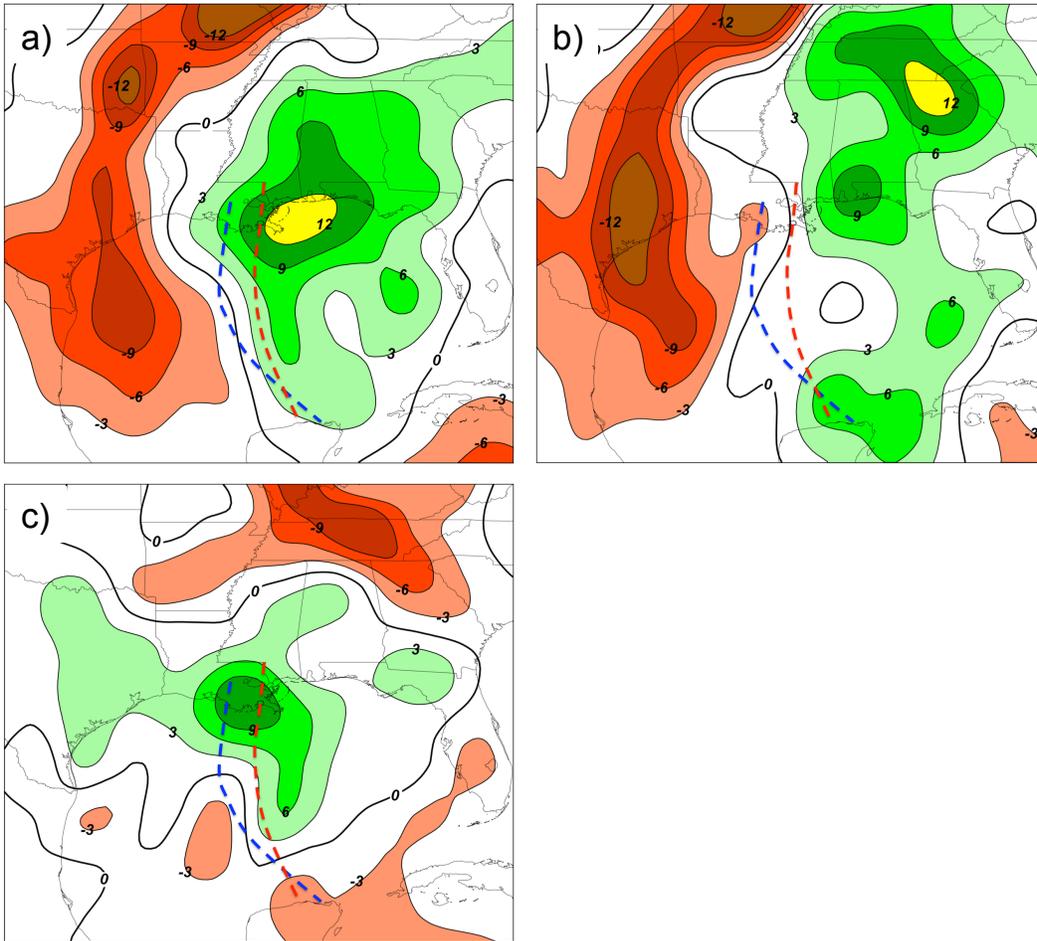


FIG. 9. Change in the magnitude of the 925 hPa (a) total, (b) geostrophic, and (c) ageostrophic poleward moisture flux over the Southeast U.S. during the 24-h period from 0000 UTC 1 May to 0000 UTC 2 May. Changes in the moisture flux greater than (less than) 3 ( $-3$ )  $\text{cm s}^{-1}$  are shaded in the green (red/brown) fill pattern every 3  $\text{cm s}^{-1}$  with 0  $\text{cm s}^{-1}$  contoured in black. Blue (red) dashed line represents the axis of maximum poleward moisture flux at 0000 UTC 1 May (2 May), as indicated in Fig. 7

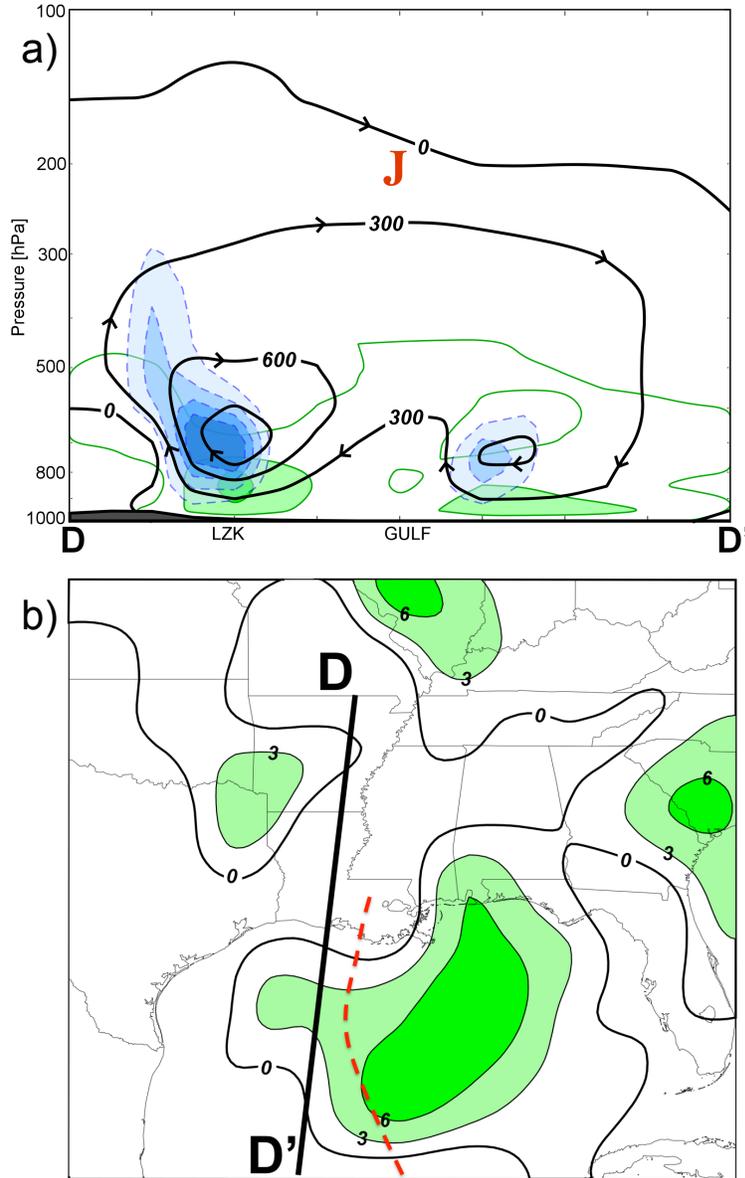


FIG. 10. (a) Cross section along the line D-D', in Fig. 7a, at 0000 UTC 1 May of Sawyer-Eliassen streamfunction every  $300 \text{ m hPa s}^{-1}$  (black lines), moisture flux associated with the Sawyer-Eliassen circulation every  $3 \text{ cm s}^{-1}$  beginning at  $0 \text{ cm s}^{-1}$  ( $0 \text{ cm s}^{-1}$  is contoured in green with the green fill pattern used for values greater than  $3 \text{ cm s}^{-1}$ ), and negative omega associated with the Sawyer-Eliassen circulation every  $1 \text{ dPa s}^{-1}$  beginning at  $1 \text{ dPa s}^{-1}$  (blue fill pattern, dashed contours). The sense of the circulation is depicted by the arrowheads plotted on the streamfunction contours, the location of the subtropical jet core is indicated by the "J", and GULF represents the Gulf Coast. (b) 925 hPa ageostrophic poleward moisture flux every  $3 \text{ cm s}^{-1}$  beginning at  $0 \text{ cm s}^{-1}$  ( $0 \text{ cm s}^{-1}$  is contoured in black with the green fill pattern used for values greater than  $3 \text{ cm s}^{-1}$ ) and the axis of maximum poleward moisture flux (red dashed line previously indicated on Fig. 7a) at 0000 UTC 1 May.

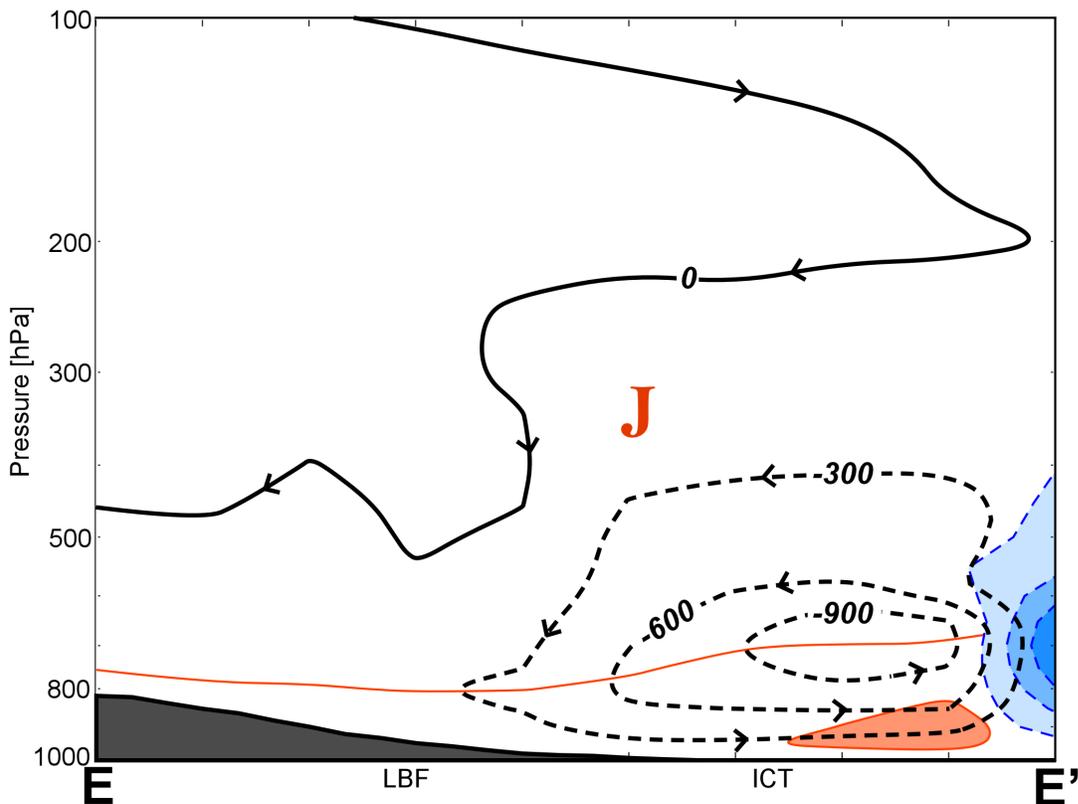


FIG. 11. Cross section along the line E-E', in Fig. 7a, at 0000 UTC 1 May of Sawyer-Eliassen streamfunction (black contours, dashed contours represent negative values) every 300 m hPa s<sup>-1</sup>, moisture flux due to the Sawyer-Eliassen circulation every -3 cm s<sup>-1</sup> beginning at 0 cm s<sup>-1</sup> (0 cm s<sup>-1</sup> is contoured in orange with the orange fill pattern used for values less than -3 cm s<sup>-1</sup>), and negative omega associated with the Sawyer-Eliassen circulation every 1 dPa s<sup>-1</sup> (blue fill pattern, dashed contours) beginning at 1 dPa s<sup>-1</sup>. The sense of the circulation is denoted by the arrowheads plotted on the streamfunction contours and the location of the polar jet core is indicated by the "J".

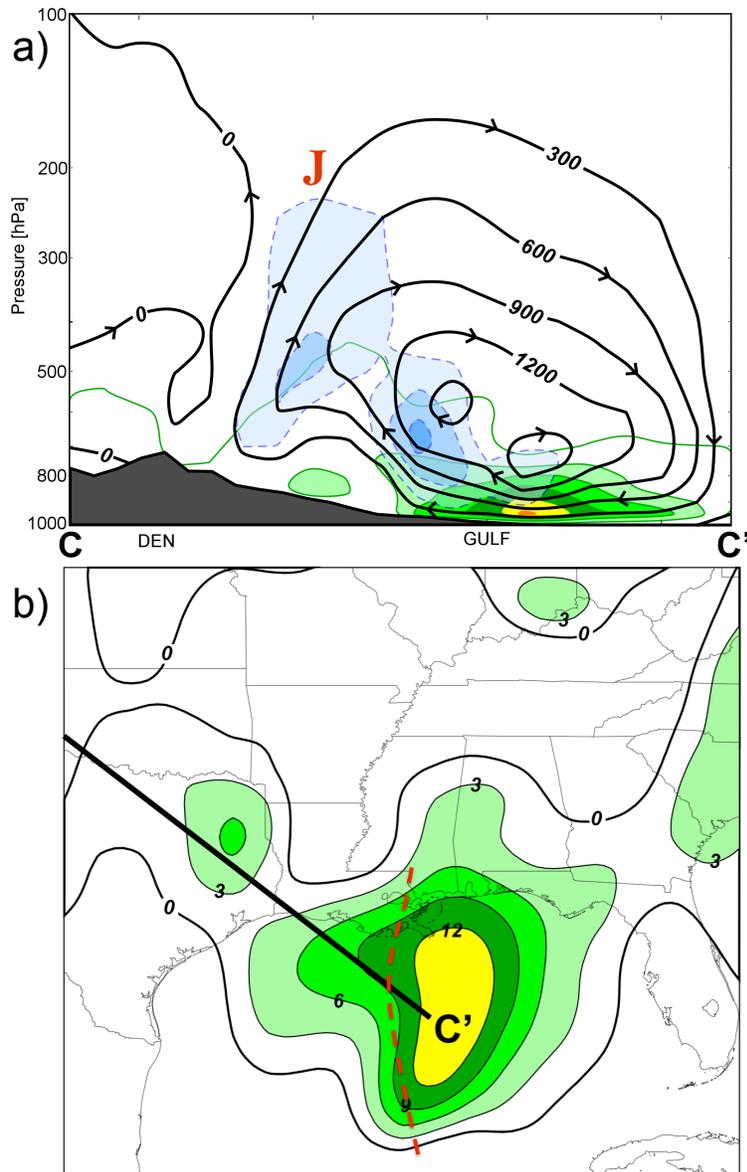


FIG. 12. (a) Cross section of Sawyer-Eliassen streamfunction along the line C-C', in Fig. 7b, at 0000 UTC 2 May. Labeling conventions are identical to those in Fig. 10a, with the "J" representing the superposed jet core. (b) As in Fig. 10b but valid at 0000 UTC 2 May.

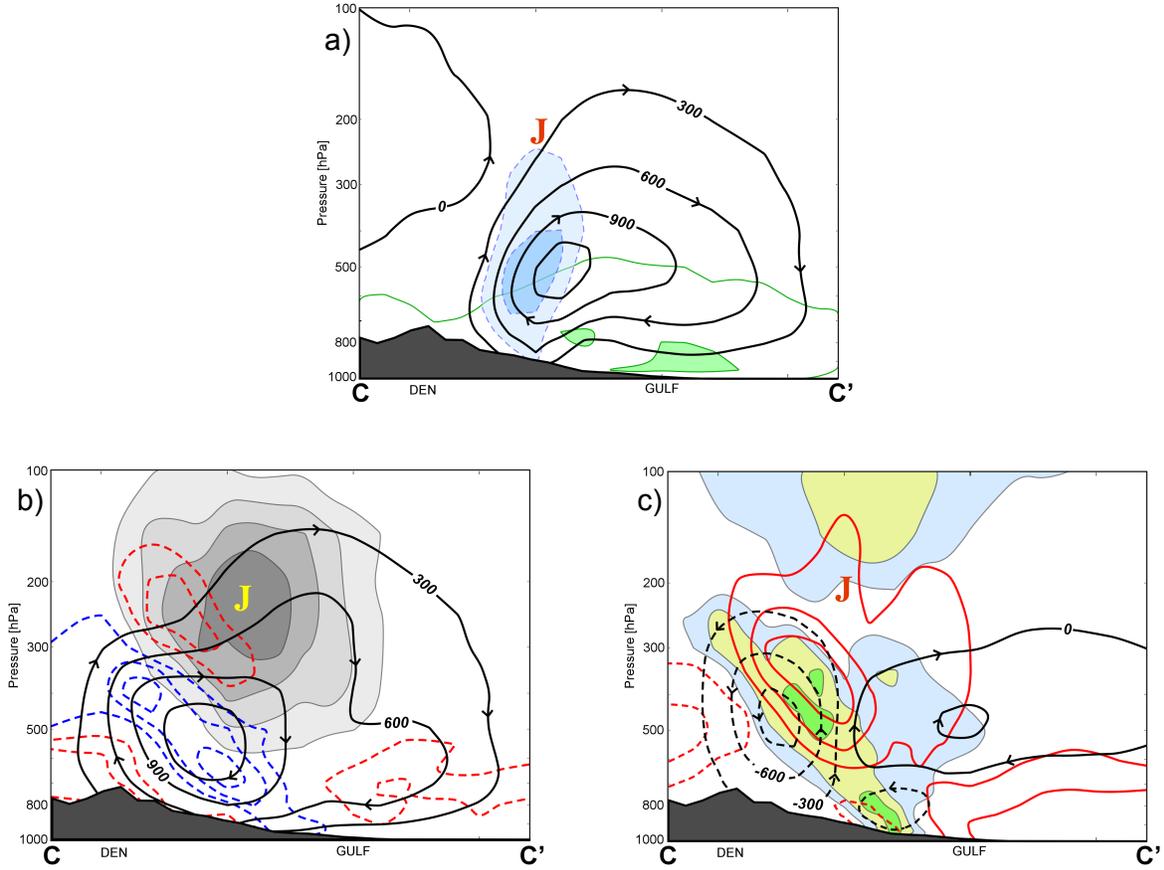


FIG. 13. Cross section along C-C', in Fig. 7b, at 0000 UTC 2 May showing: (a) the Sawyer-Eliassen streamfunction, poleward moisture flux, and negative omega (same conventions as Fig. 10a) associated with the  $Q_g$  forcing, (b) the Sawyer-Eliassen streamfunction associated with the  $Q_{SH}$  forcing (same conventions as Fig. 11), isotachs of the cross-section normal geostrophic wind (gray fill pattern) every  $10 \text{ m s}^{-1}$  beginning at  $30 \text{ m s}^{-1}$ , and the cross-section normal temperature gradient (negative, red dashed contours; positive, blue dashed contours) every  $5 \times 10^{-6} \text{ K m}^{-1}$  (zero line omitted), (c) the Sawyer-Eliassen streamfunction associated with the  $Q_{ST}$  forcing (same conventions as Fig. 11), isotachs of the along-cross section geostrophic wind with positive values oriented towards C (positive, thick red lines; negative, dashed red lines) every  $5 \text{ m s}^{-1}$  (zero line omitted), and magnitude of the along-cross section potential temperature gradient every  $10 \times 10^{-6} \text{ K m}^{-1}$  beginning at  $10 \times 10^{-6} \text{ K m}^{-1}$  (fill pattern). The “J” represents the location of the superposed jet core in all panels.

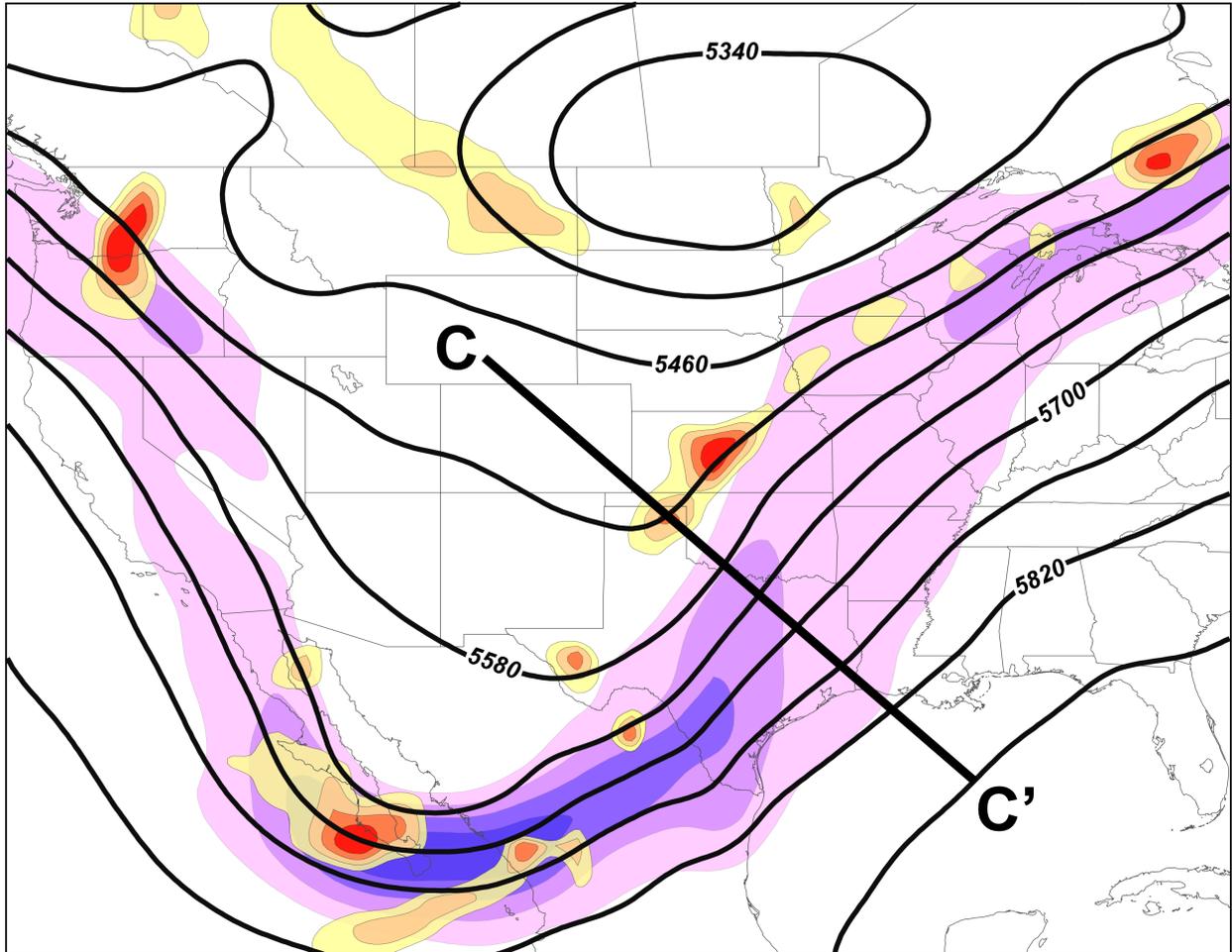


FIG. 14. 500 hPa GFS analysis at 0000 UTC 2 May with geopotential height contoured in black every 60 m, isotachs of the geostrophic wind (purple fill pattern) every  $10 \text{ m s}^{-1}$  beginning at  $30 \text{ m s}^{-1}$ , and horizontal geostrophic frontogenesis (warm colored fill pattern) every  $0.4 \text{ K (100 km)}^{-1} (3 \text{ h})^{-1}$  beginning at  $0.4 \text{ K (100 km)}^{-1} (3 \text{ h})^{-1}$ .

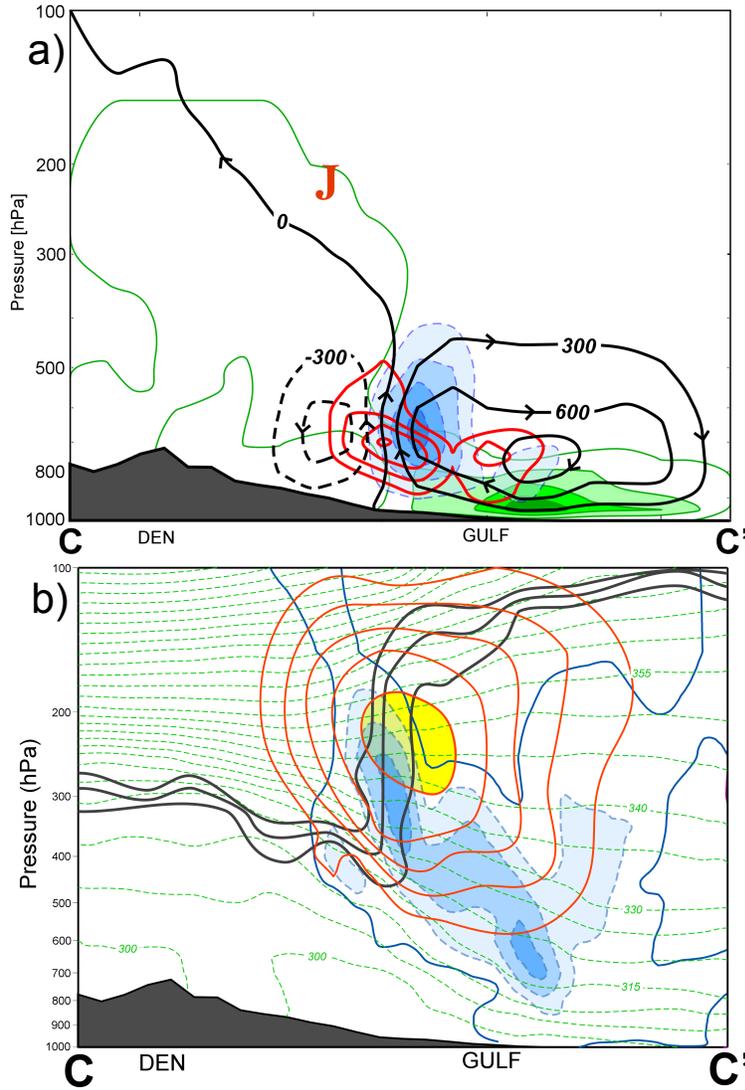


FIG. 15. (a) Sawyer-Eliassen streamfunction, poleward moisture flux, and negative omega, labeled, contoured, and shaded as in Fig. 10a, associated with the diabatic forcing. Heating labeled in  $\text{K s}^{-1}$  and contoured every  $200 \times 10^{-6} \text{ K s}^{-1}$  beginning at  $200 \times 10^{-6} \text{ K s}^{-1}$  (red contours). The “J” denotes the location of the superposed jet core. (b) 300 hPa isotachs (red contours) every  $10 \text{ m s}^{-1}$  beginning at  $30 \text{ m s}^{-1}$  with the jet core shaded yellow, 1,2,3 PVU surfaces (black contours), potential temperature every 5K (dashed green contours), and negative omega every  $2 \text{ dPa s}^{-1}$  beginning at  $0 \text{ dPa s}^{-1}$  ( $0 \text{ dPa s}^{-1}$  is contoured in blue with values greater than  $2 \text{ dPa s}^{-1}$  shaded with the blue fill pattern) from the GFS analysis at 0000 UTC 2 May 2010 along the cross section C-C’, in Fig. 7b.

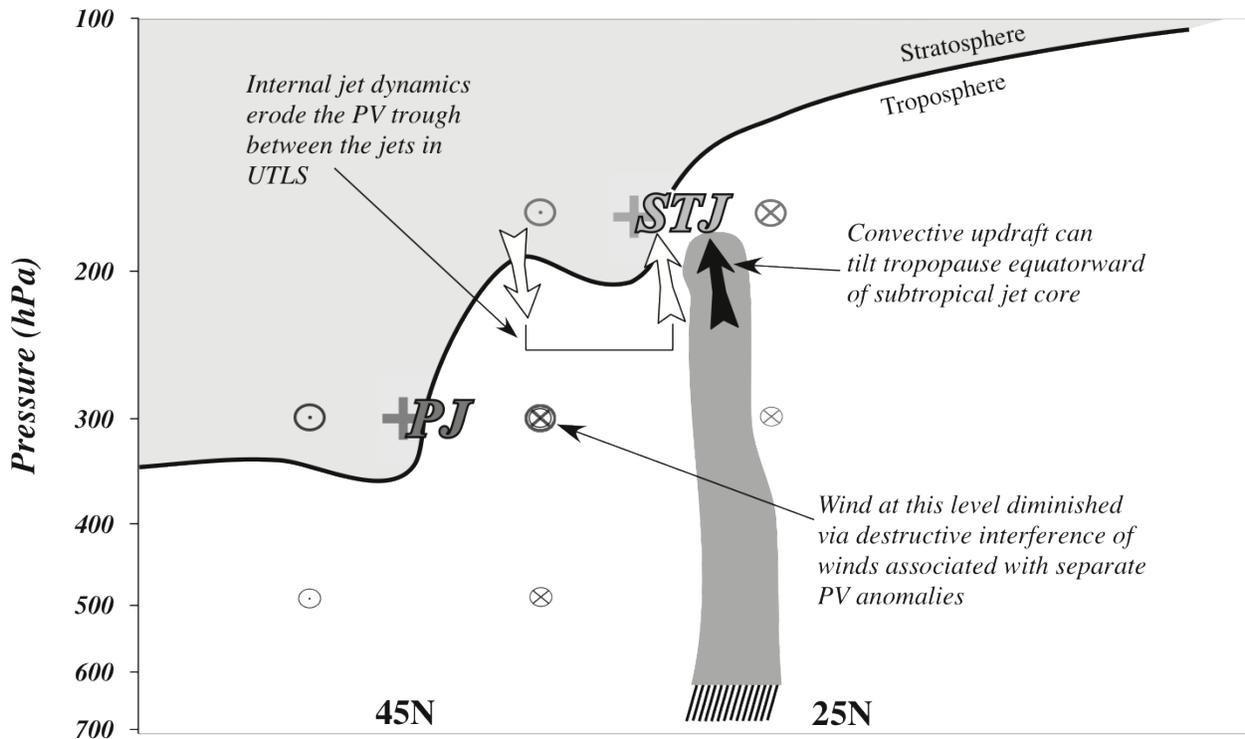


FIG. 16. Schematic vertical cross section illustrating the dynamical processes that may facilitate a superposition of the polar (PJ) and subtropical (STJ) jet. Each jet is associated with a tropopause level positive PV perturbation (signified by the + signs). Corresponding circulations at and below each perturbation are indicated by a circled  $\times$  or  $\bullet$ . Solid black line is the 1.5 PVU isosurface with the lower stratosphere shaded gray. See text for additional explanation.